

Higgs p_T at LHC: Parton Shower vs Resummation

Physics/Detector simulation meeting

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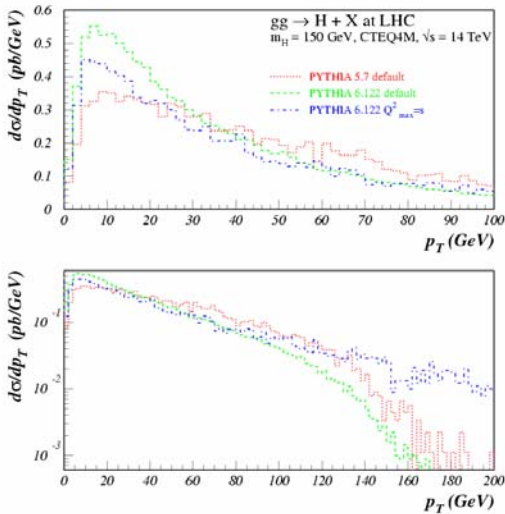
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Content

- **Motivation**
- **Parton showers**
 - **PYTHIA evolution, 5.7→6.1 and 6.1→6.2**
- **Resummation**
- **Comparisons**
 - **With experiment**
 - **$m_H = 125$ GeV @ Tevatron, LHC and SLHC**
 - **$m_H = 500$ GeV @ LHC and SLHC**
- **k_T effects in Higgs production**
- **Conclusions**

- **Work done in collaboration with C. Balazs, J. Huston, S. Mrenna and M. Tonnesman (Phys.Rev.D63:014021,2001 hep-ph/0204316, hep-ph/0011122, hep-ph/0005025)**

Motivation – Change in PYTHIA



- ▶ **Large discrepancy in Higgs p_T distribution from PYTHIA 5.7 to 6.1**
 - Especially in low p_T region (dominated by soft gluons radiations)
 - Both default versions sharply fall to 0 for $p_T = m_H$ - not very natural
 - Version 6.1 can be tuned for better behavior at large p_T
 - Pythia 6.1 considerably softer, possible concerns for $H \rightarrow \gamma\gamma$
- ▶ **Which version is more reasonable?**
 - Since there is no data, find another p_T distribution prediction to compare.

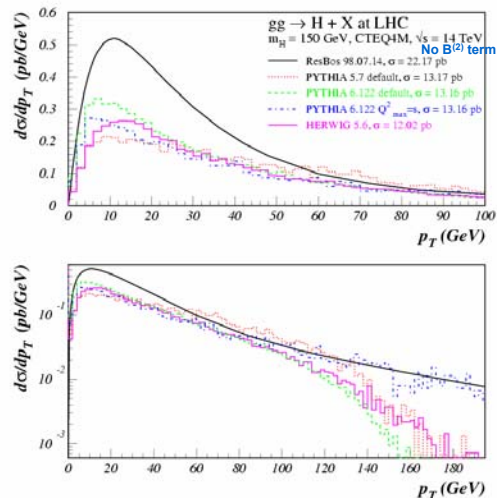
Motivation – K factor

- ▶ Monte Carlo event generators (MCEG) give a LO estimation of σ
- ▶ From the theory, in this mass region, we know:

$$K = \frac{\sigma_{NLO}}{\sigma_{LO}} \approx 1.8$$

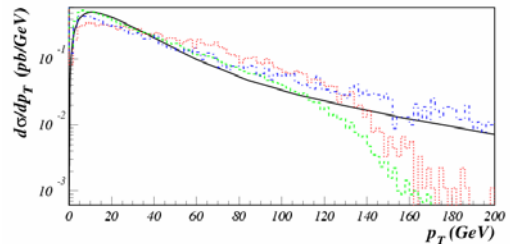
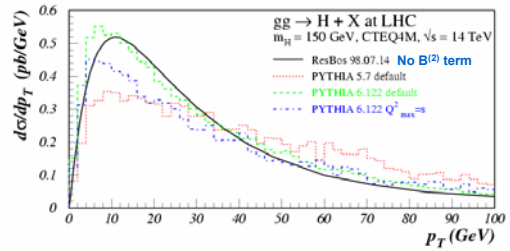
for gg fusion

- ▶ Usually, K factor implemented by scaling
 - Good kinematical agreement essential
 - Which is here obviously not the case for all MCEG versions



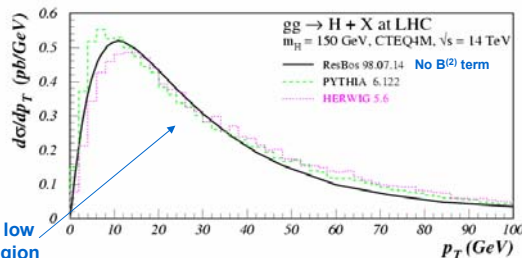
Motivation – K factor (2)

- **As a theoretical tool we take a *resummation* technique**
 - implemented in *ResBos*
 - more about it later
- **K factor implemented by scaling**
 - i.e. all cross sections are the same (that of ResBos)
- **For *inclusive* studies (f.g. $H \rightarrow ZZ^* \rightarrow 4l$) agreement between Pythia 6.1 and ResBos good enough**
- **For *exclusive* studies one should be *careful***

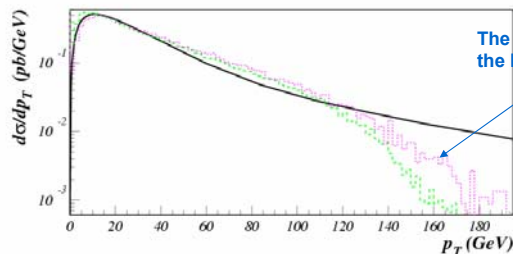


Comparison with HERWIG

- **Not so trivial, since all MCs are not equal**



Agreement good in low and intermediate region



The same as PYTHIA in the high p_T region

Parton Showers

- Monte Carlo event generators (PYTHIA, HERWIG, ISAJET etc.) based on parton showering (PSEG) are essential tools in HEP.
- PSEGs are useful because
 - they accurately describe the emission of multiple soft gluons,
 - allow a direct connection with non-perturbative models of hadronization.
- Energy-momentum is conserved at every step
 - Realistic predictions can be made for arbitrary physical quantities.
- However, the prediction of the total cross section is only accurate to LO.
- PSEGs do not accurately describe kinematical configurations with hard partons emitted at a large angle.
- Higgs p_T depends primarily on the details of the soft gluon emission from the initial state partons.
- For technical reasons, the initial state parton shower proceeds by a backward evolution.
 - Starting at the large (negative) Q^2 scale of the hard scattering.
 - And then considering emissions at lower and lower (negative) virtualities, corresponding to earlier points on the cascade, until the factorization scale is reached.
 - The transverse momentum of the initial state is build up from the whole series of splittings and boosts.

Parton Showers (2)

- The showering process is fairly independent of the hard scattering process being considered.
- Depends only on the initial state partons and the hard scale of the process.
- Parton showers in PYTHIA
 - Obey a strict ordering in virtuality: a parton that initiates a hard scattering has a larger virtuality than any other parton in the shower.
- Parton showers in HERWIG
 - Proceed via a coherent branching process in which a strict angular ordering is imposed on sequential gluon emissions.
- Because of the demonstrated importance of coherence effects, PYTHIA includes an additional veto on showers which are not also angular-ordered
 - But this does not make two schemes equivalent – some late emissions in a HERWIG shower can have virtuality larger than previous emissions.

PYTHIA Evolution: 5.7 → 6.1

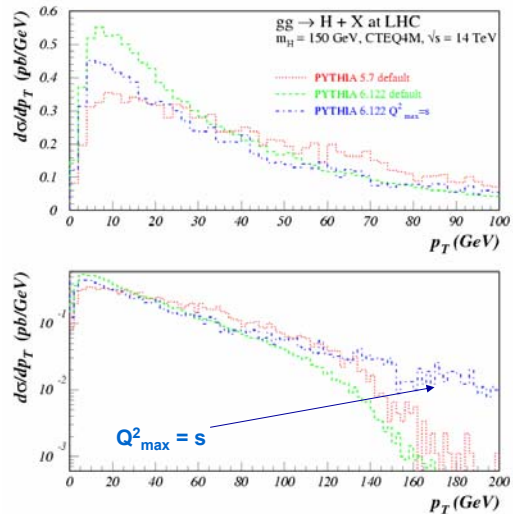
Two changes

- A cut has been placed on the combination of z and Q^2 values in a branching: $\hat{u} = Q^2 - \hat{s}(1-z) < 0$ where \hat{s} refers to the subsystem of hard scattering plus shower partons
 - corner of emissions that do not respect this requirement occurs when Q^2 value of space-like emitting parton is little changed and z value of branching is close to unity
 - necessary if matrix element corrections are to be made to process
 - net result is substantial reduction in amount of gluon radiation**
 - In principle affects all processes; in practice only gg initial states
- Parameter for minimum gluon energy emitted in space-like showers is modified by extra factor corresponding to $1/\gamma$ factor for boost to hard subprocess frame
 - result is increase in gluon radiation**
- First effect is dominant, resulting in reduction of radiation and therefore softer p_T spectrum

The above are choices, not bugs; which version is more correct?

- Compare to ResBos

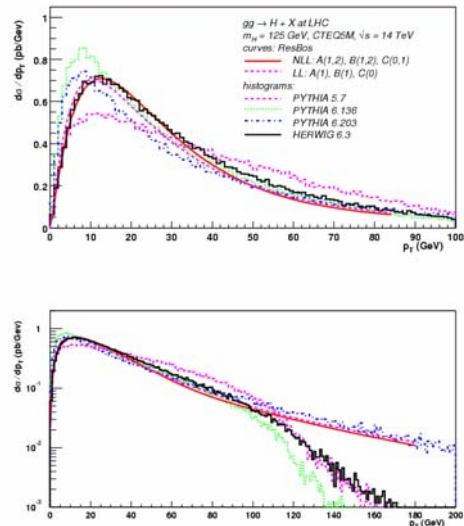
Possible implication for $H \rightarrow \gamma\gamma$ investigated and no significant effect found (K. Lassila)



PYTHIA Evolution: 6.1 → 6.2

An additional hard matrix element correction was applied to the parton shower.

- Maximum virtuality of the shower is increased from its nominal value at $Q^2_{max} = m_H^2$ up to the largest kinematically allowed value $Q^2_{max} = s$.
 - Each parton emission is corrected by ratio of the exact matrix element squared at NLO to the approximate matrix element squared given by the parton shower approximation.
- Showers generated by PYTHIA 6.2 have better shape in the high p_T region.
- Herwig still have a cutoff set by m_H .



Resummation - Introduction

- At hadron colliders, the partonic cross section for heavy boson production can receive substantial corrections at higher orders in α_s
 - This affects not only the total production rate, but also the kinematics of the heavy boson.
- At LO, the heavy boson has a $\delta(p_T^2)$ distribution in p_T^2
- At NLO
 - Real emission of a single gluon generates a contribution to $d\sigma/dp_T^2$ that behaves as $\alpha_s(p_T^2)/p_T^2$ and $\alpha_s(p_T^2)\ln(Q^2/p_T^2)/p_T^2$
 - Soft and virtual corrections are proportional to $-\delta(p_T^2)$
- At higher orders, the most singular terms follow the pattern

$$\alpha_s^n(p_T^2) \sum_{m=0}^{2n-1} \ln^m(Q^2/p_T^2) = \alpha_s^n L \equiv V^n$$

- The logarithms arise from the incomplete cancellation of the virtual and real QCD corrections
- This cancellation becomes complete for the integrated spectrum, where the real gluon can become arbitrarily soft and/or collinear to other partons
- The pattern of singular terms suggest that perturbation theory should be performed in powers of V^n instead of α_s^n
- This reorganisation of perturbative series is called **resummation**

Resummation - ResBos

- Here we use so called Collins-Soper-Sterman (CSS) formalism for p_T resummation.
 - Implemented in the ResBos program (<http://www.pa.msu.edu/~balazs/ResBos/>)
- CSS resums all of the important logarithms
 - This is achieved after a Fourier transformation with respect to p_T in the transverse coordinate b , so that the series involving the delta function and terms V_n simplifies to the form of an exponential
- The most important factors are A, B and C coefficients, which are free of large logarithmic corrections and safely calculable perturbatively as expansions in the strong coupling α_s

$$\{A, B\} = \sum_{n=1}^{\infty} \left(\frac{\alpha_s(\mu)}{\pi} \right)^n \{A^{(n)}, B^{(n)}\}, \quad C_{ij} = \sum_{n=0}^{\infty} \left(\frac{\alpha_s(\mu)}{\pi} \right)^n C^{(n)}$$

$ij = gg, gq, qq$

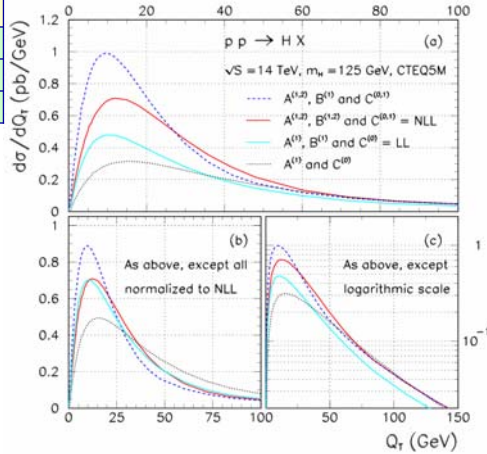
- $A^{(n)}$ coefficients are universal (i.e. process independent)
- $B^{(n)}$ coefficients are process dependant, with the exception of $B^{(1)}$
- $A^{(1)}$, $B^{(1)}$ and (approximately) $A^{(2)}$ are effectively in Monte Carlo
 - Especially HERWIG
- $A^{(1)}$, $A^{(2)}$, $B^{(1)}$ and $B^{(2)}$ for Higgs production are in current version of ResBos
 - As are $C^{(0)}$ and $C^{(1)}$ which control the NLO normalization

Resummation – LL, NLL, ...

- Two possible classifications or logarithms:
 - According to their appearance in the *Sudakov* exponent
 - Organize logs in the expanded perturbative *series*

	Sudakov	Series
LL	$A^{(1)}, C^{(0)}$	$A^{(1)}, B^{(1)}, C^{(0)}$
NLL	$A^{(1,2)}, B^{(1)}, C^{(0)}$	$A^{(1,2)}, B^{(1,2)}, C^{(0,1)}$
NNLL	$A^{(1,2,3)}, B^{(1,2)}, C^{(0,1)}$	$A^{(1,2,3)}, B^{(1,2,3)}, C^{(0,1,2)}$

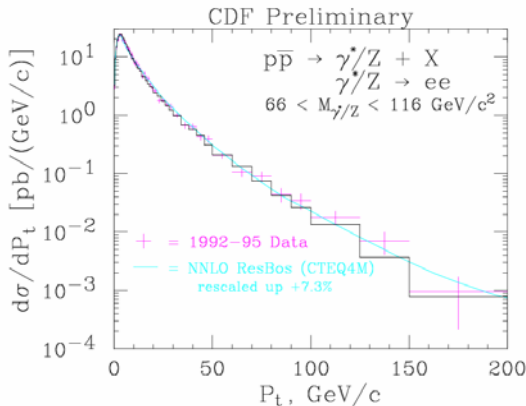
- Here we choose classification according to the perturbative expansion (*series*)
- The change from LL to NLL results in a
 - Increase of the cross section
 - Slight shift of the peak toward higher Q_T
 - But when normalizing to the same cross section, the shapes look very similar
- Effect of B(2) is surprisingly large
 - f.g. it is small for Drell-Yan production
- ResBos switches to NLO Higgs+jet matrix element at high p_T



Comparison with Experiment

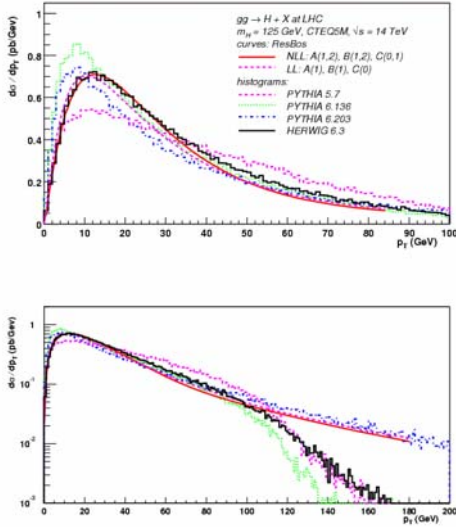
Z boson production at Tevatron

- One of the highest precision testing grounds for the effects of multiple soft-gluon emission
 - Fully differential fixed order cross section has been calculated up to $O(\alpha_s^2)$
 - $A^{(1,2)}, B^{(1,2)}, C^{(0,1)}$ coefficients are known and have been numerically implemented



- Comparison involves (only shapes are compared)
 - $Z p_T$ distribution from CDF for Run 1
 - Prediction from ResBos (rescaled upwards by 8.4%)
 - Prediction from PYTHIA 6.125 (rescaled upwards by 40%)
- PYTHIA and ResBos both describes the data well
 - Agreement of PYTHIA in the high p_T region is made possible by explicit matrix element corrections
- Low p_T region in more details later

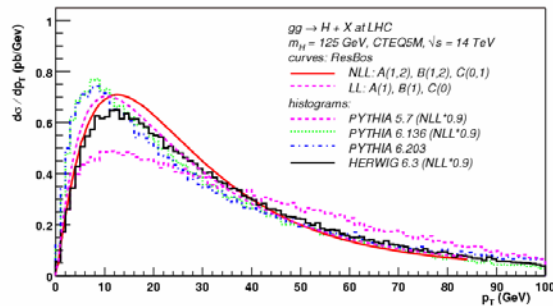
Comparison – 125 GeV Higgs at LHC



- All distributions are normalized to the ResBos NLL cross section
 - Without this the MC curves would be about a factor 2 lower than ResBos
- ResBos curves appear close to the HERWIG predictions, and somewhat less close to the predictions of PYTHIA
- In the high p_T region
 - Correct description by PYTHIA 6.2 and ResBos, because both use the exact matrix element for Higgs + jet
 - Pure parton showering description (HERWIG and old PYTHIAs) is inadequate

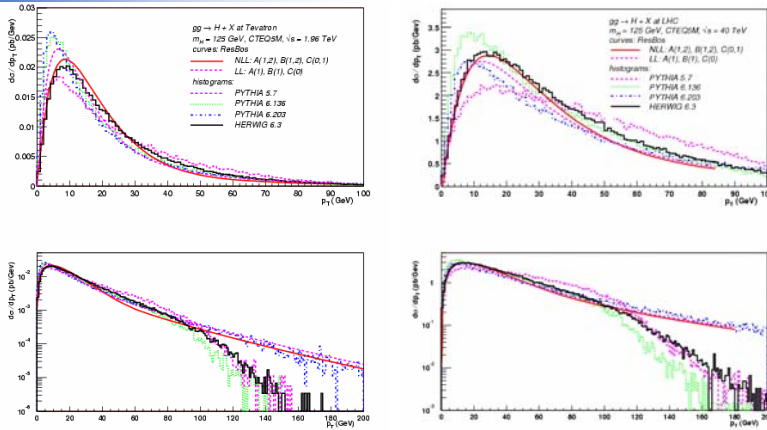
Comparison – 125 GeV Higgs at LHC (2)

- Normalizing the cross sections in the low p_T region alone



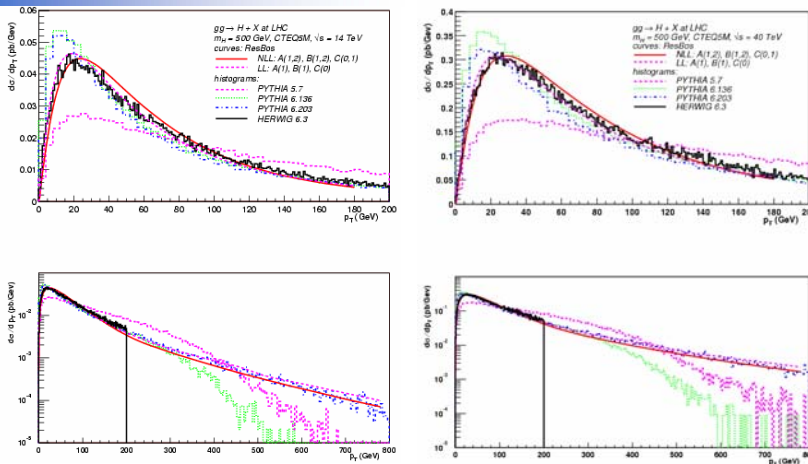
- HERWIG in very good agreement with ResBos
- PYTHIA 6.x prediction peak at lower p_T values than HERWIG and ResBos

Comparison – 125 GeV Higgs at Tevatron and SLHC



- **Average Higgs p_T increases with increasing cms energy**
 - due to the increasing phase space available for gluon emission
- **The same conclusion still valid:**
 - for low p_T good agreement ResBos-HERWIG and to less extent ResBos-PYTHIA 6.x,
 - for high p_T good agreement ResBos-PYTHIA6.2 only

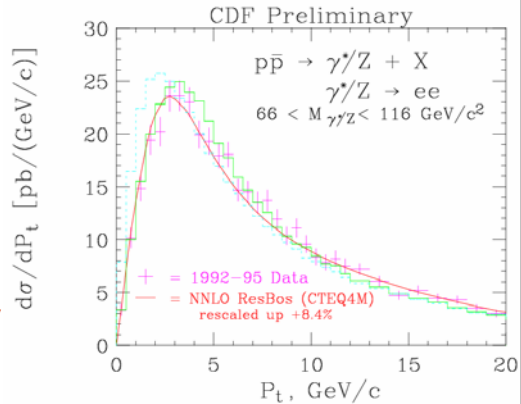
Comparison – 500 GeV Higgs at LHC and SLHC



- **Average p_T is noticeably larger than that for a 125 GeV Higgs, since the hard scale for the process is mH**
- **Conclusions for comparisons ResBos-HERWIG-PYTHIA still hold**

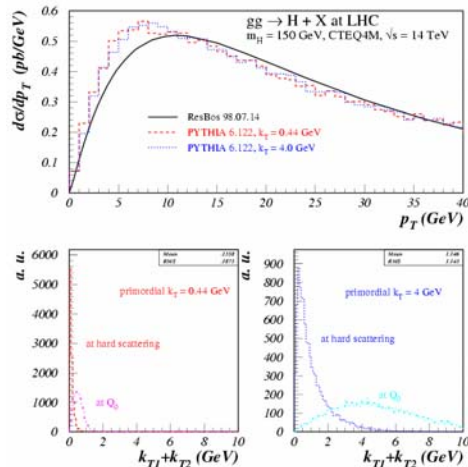
k_T Effects - Motivations

- > At a point in its evolution (typically corresponding to the virtuality of a few GeV) the parton shower is cut off and the effects on gluon emission at softer scales must be parametrized and inserted by hand
- > The parametrisation is typically expressed in a Gaussian form, in terms of non-perturbative $\langle k_T \rangle$
 - depending on the particular kinematics and initial state being investigated
- > Testing ground:
 - Z production at Tevatron**
- > PYTHIA predictions:
 - $k_T^{rms} = 0.44$ GeV (default), dashed hist.
 - $k_T^{rms} = 2.15$ GeV, solid histogram
 - Latter value was found to give the best agreement with data
- > All the difference for these two k_T s are in the low p_T region
- > High p_T region is unaffected
 - due to the number of gluon branches before collision, that carry off a sizable fraction of the original non-pert. k_T
- > Equivalent k_T in ResBos is about 2.5 GeV
- > What are k_T effects in Higgs production?



k_T Effects – Results

- > Answer: (surprisingly), no difference is observed between the predictions with the different values of k_T
- > Lower plots show the sum of the non-perturbative partonic initial state k_T 's at Q_0 and at the hard scatter scale Q
- > Most of the k_T is radiated away
 - with the effects being larger at the LHC
- > The large gluon radiation probability from a gg initial state and the greater phase space available at the LHC lead to a stronger degradation of the non-perturbative k_T than was observed with Z production at the Tevatron



Conclusions

- **Comparing the recent versions of MC event generators using the parton showers (HERWIG and PYTHIA) with the resummation calculations implemented in ResBos we come to following conclusions:**
 - **PSEG and ResBos agree well in the low and intermediate p_T region**
 - **With slightly better agreement with HERWIG than with PYTHIA**
 - **The agreement is better with increasing mass and cms energy**
 - **In high p_T region matrix element corrections are necessary for PSEGs (as implemented in recent PYTHIA)**
- **p_T spectrum softening observed from PYTHIA 5.7 \rightarrow 6.1 does not significantly affect event vertex finding in $H \rightarrow \gamma\gamma$**
- **Non-perturbative k_T effects are small in case of Higgs production**