

Nondegenerate equilibrium in (2-propyl)cycloalkyl cations. Comparison of (2-propyl)cyclopentyl and (2-propyl)cyclohexyl cation using ^{13}C NMR spectroscopy, equilibrium isotope effects and quantum chemical calculations†

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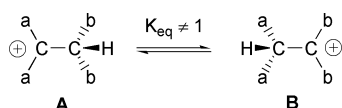
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Tertiary (2-propyl)cyclopentyl and (2-propyl)cyclohexyl carbocations were investigated by ^{13}C NMR spectroscopy in superacid solution. Both ions undergo fast nondegenerate 1,2-hydride shifts to the corresponding 2-cycloalkyl-2-propyl cations. The direction of these equilibria depends on the size of the ring. The more stable isomer of the (2-propyl)cyclopentyl cation has the formal positive charge at the endocyclic carbon atom, while the more stable isomer of the (2-propyl)cyclohexyl cation has the formal charge at the exocyclic carbon atom. The dynamic NMR results were confirmed by NMR spectroscopic measurement of the equilibrium isotope effects and rationalized by quantum chemical calculations.

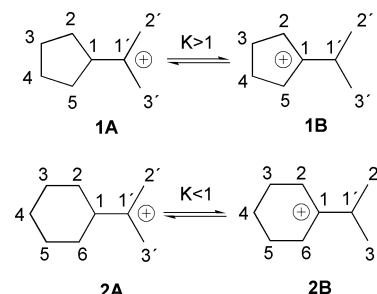
Introduction

Carbocations often show facile rearrangement reactions due to their shallow energy hypersurface. Typical examples for rapid nondegenerate rearrangements are 1,2-hydride shifts in tertiary carbocations $\text{A} \rightleftharpoons \text{B}$, such as in 2,3-dimethyl-3-pentyl cation, 2,3-dimethyl-2-hexyl cation¹ and in 2-aryl-3-methyl-2-butyl cations.²



Carbocations undergoing reversible rearrangements, which are fast relative to the NMR time scale, show NMR spectra with concentration weighted averaged chemical shifts for the interchanging sites. For a nondegenerate 1,2-hydride shift, such as $\text{A} \rightleftharpoons \text{B}$, the equilibrium constant $K_{\text{eq}} = [\text{B}]/[\text{A}] \neq 1$ and two peaks are observed for all sites which are averaged by the fast hydride shift. These peaks show temperature dependent averaged chemical shifts. The peak for the carbocation center is shifted upfield and the peak for the β -carbon is shifted downfield compared to that of static carbocations.¹ In some carbocations nondegenerate equilibria have been detected by ^1H - and ^{13}C -NMR spectroscopic methods, but often no details of those rearrangements are known.

We have investigated the 1,2-hydride shift equilibrium $\text{1A} \rightleftharpoons \text{1B}$ of the 2-cyclopentyl-2-propyl cation **1A** and the 1-(2-propyl)-cyclopentyl cation **1B** and the analogous rearrangement reaction $\text{2A} \rightleftharpoons \text{2B}$ which equilibrates the 2-cyclohexyl-2-propyl cation **2A** and the 1-(2-propyl)cyclohexyl cation **2B** (Scheme 1)



Scheme 1

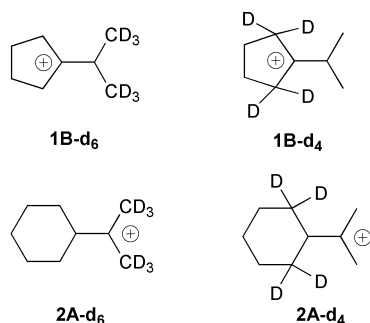
using high field ^{13}C NMR spectroscopy, equilibrium isotope effects and quantum chemical calculations. Earlier NMR studies have indicated that the equilibrium $\text{1A} \rightleftharpoons \text{1B}$ is moved toward the structure **1B** in which the positive charge is located in the ring. On the contrary, for the equilibrium $\text{2A} \rightleftharpoons \text{2B}$ the structure **2A** with the positive charge on the exocyclic carbon was found to be more stable.^{3,4}

The results we report in this study explain the different stabilities of alkyl-substituted five- and six-membered cycloalkyl cations **1B** and **2B** compared to the isomeric cycloalkyl-substituted 2-propyl cations **1A** and **2A** and give additional information on the stabilization modes and the structural and conformational consequences in these type of carbocations.

Results and discussion

The carbocations **1A/1B** and **2A/2B** and their isotopomers⁵ (**1A-d₆**/**1B-d₆**, **1A-d₄**/**1B-d₄**, **2A-d₆**/**2B-d₆** and **2A-d₄**/**2B-d₄**), as well as mixtures of deuterated and non-deuterated compounds were prepared from the corresponding alcohols with excess Lewis-acid SbF_5 and a 2 : 1 mixture of SO_2F_2 - SO_2ClF as solvent, using contemporary high vacuum matrix co-condensation techniques.⁶

† Electronic supplementary information (ESI) available: redundant coordinates for structures **1A**, **1B**, **1TS**, **2A**, **2B1**, **2B2** and **2TS**; van't Hoff plot for equilibrium $\text{1A} \rightleftharpoons \text{1B}$. See <http://www.rsc.org/suppdata/p2/b1/b107756b/>



The tertiary alcohols **1A**–**OH** or **1B**–**OH**, **1B**–**d₆**–**OH**, **1A**–**d₄**–**OH**, **2A**–**OH** or **2B**–**OH**, **2B**–**d₆**–**OH** and **2A**–**d₄**–**OH**, which were used as a precursors for the carbocations, were prepared by standard Grignard reaction from cyclopentyl- or cyclohexyl-chloride and acetone (**1A**–**OH** and **2A**–**OH**) or acetone–**d₆** (**1A**–**d₆**–**OH**, **2A**–**d₆**–**OH**), and from β,β' –**d₄**–cyclopentanone or β,β' –**d₄**–cyclohexanone and 2-chloropropane (**1B**–**d₄**–**OH**, **2B**–**d₄**–**OH**). ¹³C NMR spectra of the solutions of the carbocations were taken at different temperatures and assigned using standard techniques. Some representative spectra are shown in Fig. 1.

Temperature dependent ¹³C NMR chemical shifts for **1A**/**1B** and **2A**/**2B** are given in Table 1.

The chemical shifts of the averaged signal of the interchanging sites depend on the equilibrium constant. The relative population of the two species **A** and **B** is given by the equilibrium constant $K_{eq} = [B]/[A]$. The averaged chemical shifts δ_1 and $\delta_{1'}$ of the two peaks of a pair of carbons such as C1 and C1' which are interchanged by the fast 1,2-hydride shifts are given in eqn. (1) and (2),

$$\delta_1 = \frac{\delta_1(A) + K_{eq} \delta_1(B)}{1 + K_{eq}} \quad (1)$$

$$\delta_{1'} = \frac{\delta_{1'}(A) + K_{eq} \delta_{1'}(B)}{1 + K_{eq}} \quad (2)$$

where $\delta_1(A)$ and $\delta_{1'}(A)$ represent the chemical shifts of carbons C1 and C1' in the structure **A**, and $\delta_1(B)$ and $\delta_{1'}(B)$ represent the chemical shifts of the same carbons in the structure **B**. The equilibrium constant K_{eq} is obtained from the chemical shift difference $\Delta = \delta_1 - \delta_{1'}$ for the two carbons according to eqn. (3).

$$K_{eq} = \frac{\delta_1(A) - \delta_{1'}(A) + \Delta}{\delta_1(B) - \delta_{1'}(B) - \Delta} \quad (3)$$

Eqn. (3) can be applied to any two interchanging sites. If the chemical shift difference Δ between the two interchanging sites is small, however, the error of the equilibrium constant will be large. If an equilibrium is degenerate, structure **A** and **B** are identical, $\delta_1(A) - \delta_{1'}(A) = \delta_1(B) - \delta_{1'}(B) = \Delta$ and eqn. (3) converts to $K_{eq} = (\Delta + \delta)/(\Delta - \delta)$.^{7,8}

For the determination of the experimental constant (K_{obs}) using eqn. (3), the ¹³C chemical shifts of the slow exchange limiting structures **A** and **B** are required. For low barrier rearrangement processes slow exchange spectra are often not accessible by experimental NMR spectroscopy in solution. The chemical shifts for these structures were therefore determined using quantum chemical calculations with the GIAO-B3LYP/6-311G(d,p) method for B3LYP/6-31G(d) optimized geometries. We have shown for a model set of alkyl carbocations by comparison of calculated chemical shifts with experimental shifts, that at this level of calculation, after applying a linear scaling method, an accuracy of $\Delta\delta < \pm 2$ ppm is obtained.⁹

Geometry optimizations and frequency calculations for the structures involved in the equilibria **1A** $\rightleftharpoons$ **1B** and **2A** $\rightleftharpoons$ **2B**,

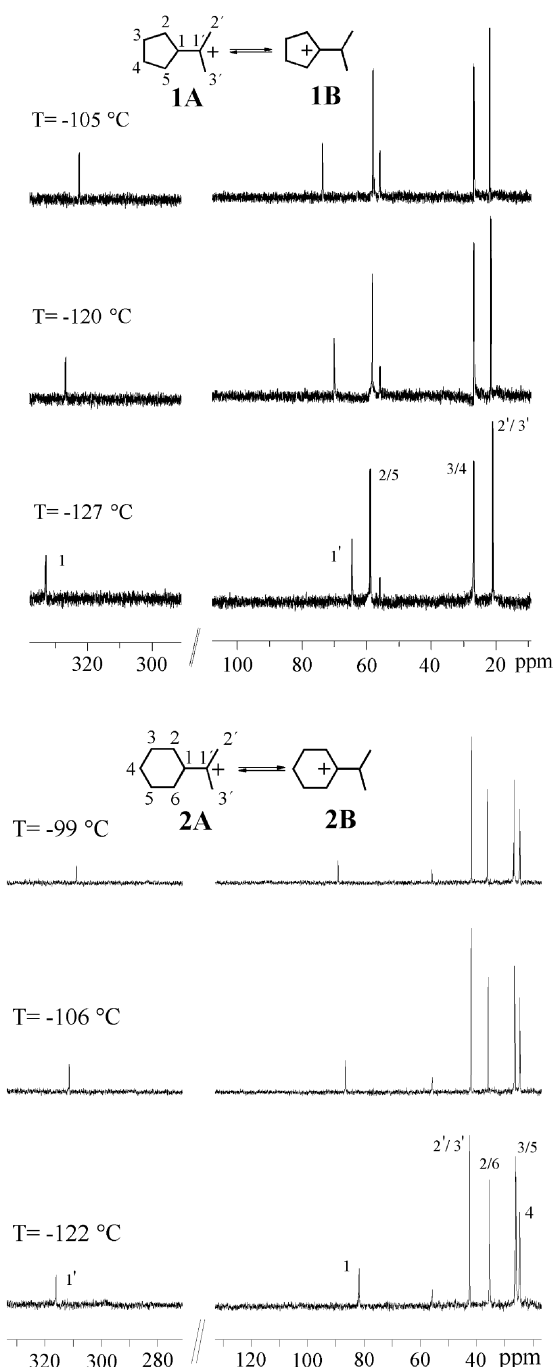


Fig. 1 ¹³C NMR spectra (100 MHz) of **1A**/**1B** and **2A**/**2B** in SO_2ClF – SO_2F_2 solution at selected temperatures.

respectively, were carried out at the B3LYP/6-31G(d) and MP2/6-31G(d) levels of theory using the Gaussian 98 program packages.¹⁰ The energies of the minimum structures of the carbocations **1A**, **1B** and **2A**, **2B1**, **2B2** are presented in Table 2. Selected bond lengths and angles are shown in Fig. 2 and 3.

The transition state structures **1TS** for the rearrangement **1A** $\rightleftharpoons$ **1B** and **2TS** for the rearrangement **2A** $\rightleftharpoons$ (**2B1** + **2B2**)/2 were located and characterized as first order saddle points (NImag = 1) at the B3LYP/6-31G(d) and the MP2/6-31G(d) level. Selected bond lengths and angles are also shown in Fig. 2 and 3, respectively.

The optimized MP2/6-31G(d) geometries were subjected to single point energy calculations at MP4 level, thus our final level is MP4(SDTQ)/6-31G(d)//MP2/6-31G(d). The calculated ¹³C chemical shifts (δ_{calc}) at the GIAO-B3LYP/6-311G(d,p)//B3LYP/6-31G(d) level of theory and the scaled chemical shifts

Table 1 Temperature dependent ^{13}C (100 MHz) NMR chemical shifts of cations **1A/1B** and **2A/2B**, and the corresponding equilibrium constants

Cation	$T/^{\circ}\text{C}$	δ (ppm)					$\Delta = \delta_1 - \delta_{1'}$ (ppm)	K_{obs}^a	
		C3,C4	C2,C5	C1	C1'	C2',C3'			
1A/1B	−127	26.73	58.61	332.73	64.24	20.76	268.50	28.97	
	−120	2.71	58.33	329.73	66.86	21.04	259.99	21.38	
	−112	26.70	58.08	326.55	69.78	21.34	256.77	16.56	
	−105	26.70	57.76	322.48	73.42	21.74	249.06	12.81	
	−97	26.67	57.38	318.27	77.21	22.06	241.06	10.24	
	−86	26.67	56.97	315.56	79.88	22.54	235.68	7.16	
2A/2B		C4	C3,C5	C2,C6	C1	C1'	C2',C3'		
	−122	24.78	25.87	35.34	81.67	316.20	42.46	−234.53	0.0131
	−115	24.76	25.91	35.55	83.89	314.01	42.22	−230.11	0.0222
	−106	24.72	25.86	35.81	86.59	311.18	41.94	−224.70	0.0336
	−99	24.69	26.02	36.04	89.04	308.65	41.67	−219.61	0.0442
	−86	24.63	26.15	36.51	94.28	303.27	41.12	−208.99	0.0676

^a $K_{\text{obs}} = (\delta_1(\mathbf{A}) - \delta_{1'}(\mathbf{A}) + \Delta)/(\delta_1(\mathbf{B}) - \delta_{1'}(\mathbf{B}) - \Delta)$ where $\delta_1(\mathbf{A})$, $\delta_{1'}(\mathbf{A})$, $\delta_1(\mathbf{B})$ and $\delta_{1'}(\mathbf{B})$ are the calculated ^{13}C NMR chemical shifts (see Table 3) of the minimum energy structures **1A**, **1B**, **2A** and **2B**, respectively.

^a $K_{\text{obs}} = (\delta_1(\text{A}) - \delta_{1'}(\text{A}) + \Delta)/(\delta_1(\text{B}) - \delta_{1'}(\text{B}) - \Delta)$ where $\delta_1(\text{A})$, $\delta_{1'}(\text{A})$, $\delta_1(\text{B})$ and $\delta_{1'}(\text{B})$ are the calculated ^{13}C NMR chemical shifts (see Table 3) of the minimum energy structures **1A**, **1B**, **2A** and **2B**, respectively.

(δ_{scaled}) for the minimum energy structures are presented in Table 3.

The 2-cyclopentyl-2-propyl cation (**1A**) $\rightleftharpoons$ 1-(2-propyl)cyclopentyl cation (**1B**) rearrangement

The five signals which are observed in the ^{13}C NMR spectrum of cation **1A/1B** (Table 1, Fig. 1) show temperature dependent chemical shifts, indicating that the cation undergoes a reversible rearrangement which is fast on the NMR time scale. The two structures which are interchanged by a nondegenerate hydride shift are the 2-cyclopentyl-2-propyl cation (**1A**) and the 1-(2-propyl)cyclopentyl cation (**1B**). The relative shielding of the methylene groups C2/C5 (58.3 ppm, −120 $^{\circ}\text{C}$, triplet) vicinal to the formal carbocation center C1 in **1B** and the shielding of the methyl groups C2'/C3' (21.0 ppm, −120 $^{\circ}\text{C}$, quartet) which are vicinal to the formal carbocation center C1' in **1A** indicate that in the equilibrium **1A** $\rightleftharpoons$ **1B** the isomer **1B** which has the formal charge located on the endocyclic C1 carbon atom is favoured. This is consistent with the observed temperature dependence of the chemical shifts (Table 1) and also in accord with earlier conclusions by Okazawa and Sorensen,³ and Saunders *et al.*⁴ At higher temperatures the population of the less stable isomer **1A** increases, affecting the averaged chemical shifts. The ring carbons C1, C2/C5 and C3/C4 move upfield and the exocyclic carbons C1' and C2'/C3' move downfield.

Both cation structures **1A** and **1B** have conformations favourable for β -CC-hyperconjugative stabilization of the positive charge (Fig. 2).¹¹ The cyclopentyl ring substituent of the 2-propyl cation **1A** has an envelope conformation, carbons C1, C2, C3 and C5 are coplanar. The C1'–C2'–C3' plane of the isopropyl group is perpendicular (87 $^{\circ}$) to C1–C2–C3–C5 thus the main axis of the formal vacant $2p\pi$ orbital at carbon C1' is eclipsed with the C1–C5 bond. Characteristic structural distortions in **1A** accompany the hyperconjugative stabilization of the positive charge formally located at the C1' carbon. The C1–C5 cyclopentyl bond ($d(\text{C1}–\text{C5}) = 1.664$ or 1.704 Å, at B3LYP/6-31G(d) and MP2/6-31G(d) level, respectively) is considerably elongated due to hyperconjugation (shown in black in Fig. 2) with the $2p\pi$ orbital at the carbon C1' as compared to the other vicinal cyclopentyl bond ($d(\text{C1}–\text{C2}) = 1.539$ or 1.520 Å; at B3LYP/6-31G(d) or MP2/6-31G(d) level of theory). Even more pronounced is the variance of the two bond angles C1'–C1–C5 (95.8 and 82.2 $^{\circ}$) and C1'–C1–C2 (123.7 and 125.7 $^{\circ}$) indicating the hyperconjugatively induced compression of the C1'–C1–C5 bond angle. The onset of bridging between C1' and C5 is indicated by a shorter distance ($d(\text{C1}'–\text{C5}) = 2.058$ Å) as compared to $d(\text{C1}'–\text{C2}) = 2.600$ Å (MP2/6-31G(d) level).

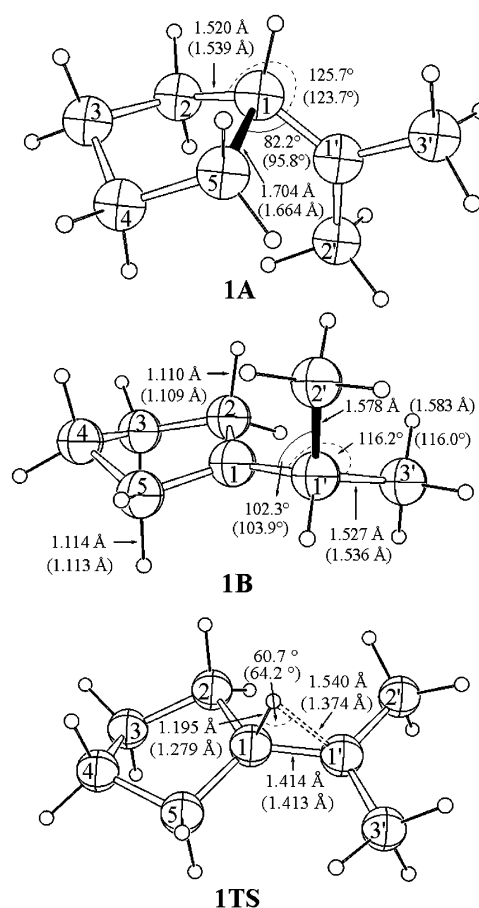


Fig. 2 MP2/6-31G(d) optimized geometries of cations **1A**, **1B** and the transition state structure **1TS** (B3LYP/6-31G(d) values in parentheses). Bonds involved in hyperconjugation are shown in black.

The cyclopentyl cation structure **1B** has a slightly twisted ring conformation. Carbon C4 is above and carbon C3 is below the plane C1–C2–C5–C1'. The β -C–C bond of one of the methyl groups (C2'–C1') forms an angle of 90 $^{\circ}$ to this plane, and is thus optimally aligned with the formally vacant $2p\pi$ orbital at C1 for hyperconjugative interaction. This bond ($d(\text{C2}'–\text{C1}') = 1.583$ or 1.578 Å) is elongated as compared to the C–C-bond to the other methyl group ($d(\text{C1}'–\text{C3}') = 1.536$ Å or 1.527 Å) and has a reduced bond angle (C2'–C1'–C1, 103.9 or 102.3 $^{\circ}$) as compared to that of the other CH₃-group (C3'–C1'–C1, 116.0 or 116.2 $^{\circ}$, B3LYP/6-31G(d) or MP2/6-31G(d),

Table 2 Total and relative energies for carbocations **1A**, **1B**, **2A**, **2B1** and **2B2** and for transition state structures **1TS** and **2TS** in vacuum ($\epsilon = 1$) and solvent ($\epsilon = 30.0$)

Cation	$E_{\text{gas}}/\text{hartree}$		$E_{\text{solv}}/\text{hartree}^a$		$\Delta E_{\text{gas}}/\text{kcal mol}^{-1}$		$\Delta E_{\text{solv}}/\text{kcal mol}^{-1}$		$E_{\text{gas}}/\text{hartree}$	$\Delta E_{\text{gas}}/\text{kcal mol}^{-1}$
	B3LYP/ 6-31G(d) (ZPE) ^b [NImag] ^c	MP2/ 6-31G(d) (ZPE) ^b [NImag] ^c	B3LYP/ 6-31G(d) (ZPE) ^b [NImag] ^c	MP2/ 6-31G(d) (ZPE) ^b [NImag] ^c	B3LYP/ 6-31G(d) + ZPE	MP2/ 6-31G(d) + ZPE	B3LYP/ 6-31G(d) + ZPE	MP2/ 6-31G(d) + ZPE	MP4/6-31G(d)// MP2/6-31G(d) (ZPE)	
1A	−313.610958 (0.213313) [0]	−312.451585 (0.218267) [0]	−313.682384 (0.213313) [0]	−312.522486 (0.218267) [0]	1.60	0	1.50	0	−312.573233 (0.218267)	0.54
1B	−313.612480 (0.212465) [0]	−312.448807 (0.216891) [0]	−313.684033 (0.212465) [0]	−312.520181 (0.216891) [0]	0	0.58	0	0.88	−312.572712 (0.216891)	0
1TS	−313.599733 (0.210580) [1]	−312.438571 (0.214935) [1]	−313.670599 (0.210580) [1]	−312.509886 (0.214935) [1]	7.25	5.82	6.82	6.08	−312.561922 (0.214935)	5.60
2A	−352.932279 (0.242476) [0]	−351.622437 (0.247539) [0]	−353.001729 (0.242476) [0]	−351.691344 (0.247539) [0]	0.29	1.34	0	4.56	−351.761354 (0.247539)	1.19
2B1	−352.934061 (0.243784) [0]	−351.625611 (0.248576) [0]	−353.001387 (0.243784) [0]	−351.699659 (0.248576) [0]	0	0	1.04	0	−351.764288 (0.248576)	0
2B2	−352.932552 (0.242469) [0]	−351.621897 (0.247558) [0]	−353.000229 (0.242469) [0]	−351.698550 (0.247558) [0]	0.12	1.69	0.94	0.06	−351.761334 (0.247558)	1.21
2TS	−352.918036 (0.240094) [1]	−351.610006 (0.245143) [1]	−352.987002 (0.240094) [1]	−351.679755 (0.245143) [1]	7.74	7.64	7.75	10.34	−351.748891 (0.245143)	7.58

^a SCRF(IPCM) model and solvent relative permittivity of $\epsilon = 30.0$ were used. B3LYP/6-31G(d) and MP2/6-31G(d) single point energy SCRF calculations were performed for B3LYP/6-31G(d) and MP2/6-31G(d) optimized geometries, respectively. ^b Zero point vibrational energy (not scaled) obtained from frequency calculation. ^c Number of imaginary frequencies obtained from frequency calculation.

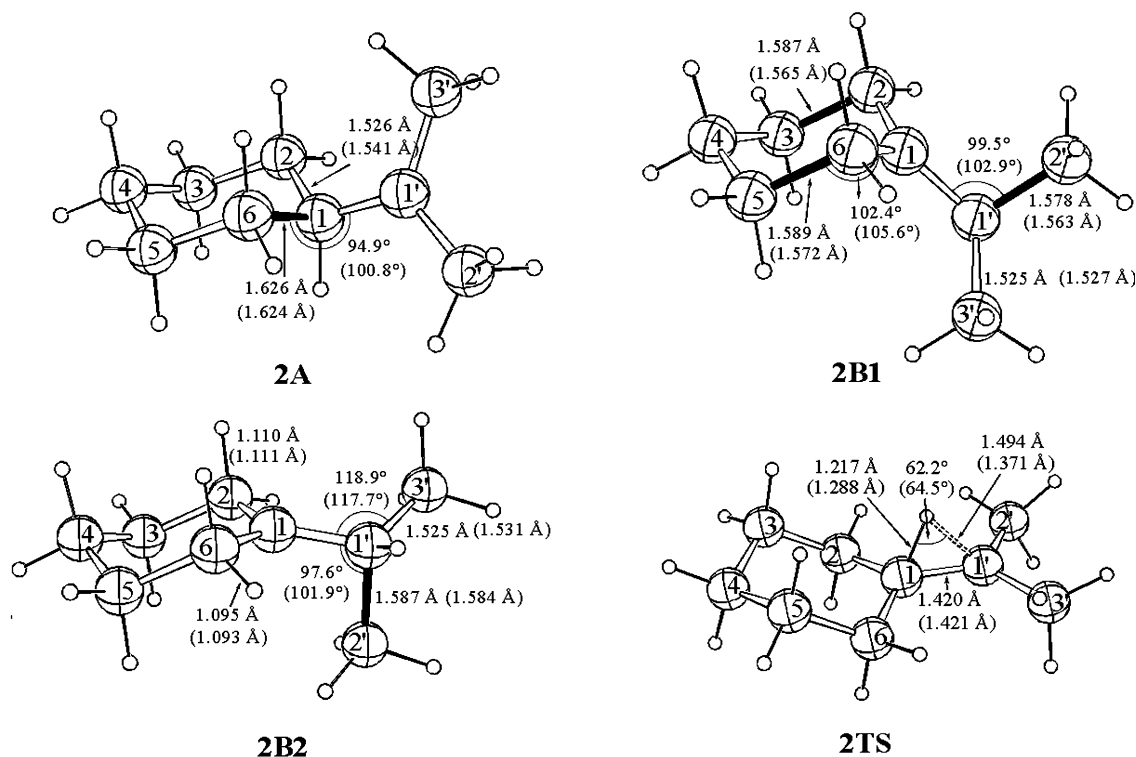


Fig. 3 MP2/6-31G(d) optimized geometries of cations **2A**, **2B1** and **2B2** and the transition state structure **2TS** (B3LYP/6-31G(d) values in parentheses). Bonds involved in hyperconjugation are shown in black.

respectively). The twisted ring conformation calculated for the 1-(2-propyl)cyclopentyl cation (**1B**) is in agreement with previous conclusions based on experimental NMR spectroscopic investigations of equilibrium isotope effects in β -CD₂-

and CHD-deuterated isotopomers of the 1-methylcyclopentyl cation¹² and is also in accord with *ab initio* calculations for the twisted conformation of the parent cyclopentyl cation.¹³ The cyclopentyl cation **1B** is calculated to be 1.6 kcal mol^{−1} (B3LYP/

Table 3 ^{13}C NMR chemical shifts of 2-cyclopentyl-2-propyl cation (**1A**), 1-(2-propyl)cyclopentyl cation (**1B**), 2-cyclohexyl-2-propyl cation (**2A**) and 1-(2-propyl)cyclohexyl cation (**2B1** and **2B2**) calculated at the GIAO-B3LYP/6-311G(d,p)//B3LYP/6-31G(d) level (δ_{calc}) and the corresponding scaled values (δ_{scaled}),^a referenced to TMS^b

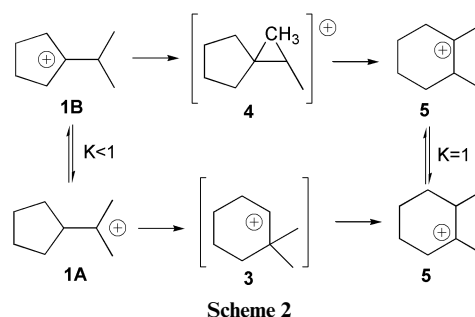
		Carbon atom					
Cation		C3,C4 ^c		C2,C5 ^c	C1	C1'	C2',C3' ^c
1A	δ_{calc}	36.22	47.00	86.85	299.25	39.98	
	δ_{scaled}	25.64	28.97	75.89	290.68	37.22	
1B	δ_{calc}	32.12	63.54	357.62	62.82	28.77	
	δ_{scaled}	26.77	56.65	348.06	56.07	17.25	
		C1'	C3,C2 ^c	C2,C6 ^c	C1	C1'	C2',C3' ^c
2A	δ_{calc}	28.17	31.02	44.05	85.99	325.30	46.52
	δ_{scaled}	27.48	29.73	26.29	75.18	316.29	42.62
2B1	δ_{calc}	32.19	60.65	59.67	333.66	65.39	24.54
	δ_{scaled}	30.80	51.55	53.47	324.51	58.18	15.03
2B2	δ_{calc}	27.17	33.02	58.91	345.56	68.68	23.63
	δ_{scaled}	26.66	27.60	52.84	336.21	60.21	14.85

^a For the scaling equation and scaling parameter see ref. 9. ^b The absolute shielding values (^{13}C : 183.86; ^1H 31.92; ^{29}Si 339.94) for tetramethylsilane (TMS; T_d symmetry) were obtained using the GIAO-B3LYP/6-311G(d,p)//B3LYP/6-31G(d) method. ^c Averaged chemical shift of the two corresponding carbon atoms.

6-31G(d)) or 0.5 kcal mol⁻¹ (MP4/6-31G(d)//MP2/6-31G(d)) more stable than the 2-propyl cation structure **1A**, which is in accord with the experimental findings. At the MP2/6-31G(d)//MP2/6-31G(d) level, however, **1B** is calculated to be 0.6 kcal mol⁻¹ less stable than **1A**. As expected for such a flat energy surface, the energy differences show some dependence on the calculation method used. The μ -hydrido-bridged transition state structure **1TS** (Fig. 2) is 7.3 kcal mol⁻¹ (B3LYP/6-31G(d)) or 5.6 kcal mol⁻¹ (MP4/6-31G(d)//MP2/6-31G(d)) higher in energy than **1B**. This energy difference is a measure for the barrier of the 1,2-hydride shift interconverting the structure **1B** and **1A**. No minimum energy structures with a conformation suitable for β -methine-C-H/C⁺-2p π -hyperconjugation was found. Initial starting structures similar to **1A** and **1B** but with the β -C1-H and β -C1'-H-bond, respectively, aligned with the formally vacant 2p π -orbital at the carbocation center, on geometry optimization, at the DFT-hybrid and MP2 level, converged to the β -CC hyperconjugation type minimum structures **1A** and **1B** respectively. Equilibrium constants K_{obs} for the equilibrium **1A** $\rightleftharpoons$ **1B** were determined from ^{13}C chemical shifts of the signals for C1 and C1' measured between -127 and -86 °C (Table 1) using eqn. (3) and the scaled calculated chemical shifts for C1 and C1' in **1A** and **1B** (Table 3). A linear van't Hoff plot ($r = 0.999$) was obtained plotting $\ln K$ vs. $1/T$. From the slope and the intercept the enthalpy $\Delta H^\circ = -1.8 \pm 0.4$ kcal mol⁻¹ and the entropy $\Delta S^\circ = 5.7 \pm 0.3$ cal K⁻¹mol⁻¹ were determined for the rearrangement reaction **1A** $\rightleftharpoons$ **1B**.

The hydride shift **1A** $\rightleftharpoons$ **1B** is fast on the NMR time scale even at the lowest temperatures (-127 °C) accessible in solution. Kinetic line broadening is observed for all but the peaks for C3,C4 at temperatures below -100 °C. The narrow temperature range (-127 to -86 °C) does not allow accurate line shape analysis, thus only the lower¹⁴ and the upper¹⁵ limits of the reaction rate at -127 °C were estimated to be $4.3 \times 10^4 < k < 9.0 \times 10^5$.¹⁶ The corresponding limits for the free energy of activation at -127 °C are $5.0 < \Delta G^\ddagger < 5.2$ kcal mol⁻¹. This is in fair agreement with the calculated energy difference between **1B** and **1TS** (7.3 and 5.6 kcal mol⁻¹ at the B3LYP/6-31G(d) and MP4/6-31G(d)//MP2/6-31G(d) level, respectively). Above -100 °C a slow but irreversible rearrangement of **1A/1B** to 1,2-dimethylcyclohexyl cation **5** was observed.¹¹ The rate con-

stant for this rearrangement process **1A/1B** $\rightarrow$ **5** at -99 °C is $k = 2.7 \times 10^{-3}$ s⁻¹, and $\Delta G^\ddagger$ (-99 °C) is 12.0 kcal mol⁻¹. The formation of carbocation **5** can be rationalized by a mechanism suggested in Scheme 2.

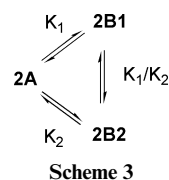


At least two different pathways are possible. Both the expansion of the five membered ring of the 2-(2-cyclopentyl)propyl cation (**1A**) and 1,2-methyl shift in the 1-(2-propyl)cyclopentyl cation (**1B**) yield the tertiary carbocation **5**, which is known to undergo a degenerate low barrier 1,2-hydride shift.¹⁵

The 2-cyclohexyl-2-propyl cation (**2A**) $\rightleftharpoons$ 1-(2-propyl)cyclohexyl cation (**2B**) rearrangement

The ^{13}C NMR spectrum of the mixture of 2-cyclohexyl-2-propyl cation (**2A**) and 1-(2-propyl)cyclohexyl cation (**2B**) shows six resonances with temperature dependent chemical shifts (Table 1, Fig. 1), proving the reversible rearrangement **2A** $\rightleftharpoons$ **2B**, which is fast on the NMR time scale. The relative shielding of the methylene groups C2,C6 (35.3 ppm, -122 °C, triplet) vicinal to the formal carbocation center C1 in **2B** and that of the methyl carbons C2',C3' (42.5 ppm, -122 °C, quartet) which are vicinal to the formal carbocation center C1' in **2A** indicate that the equilibrium **2A** $\rightleftharpoons$ **2B** is shifted towards the isomer **2A** which has a charge located on the exocyclic carbon C1'. This is consistent with the direction of the temperature dependent change of the averaged chemical shifts (Table 1). The averaged peaks for the ring carbons C1, C2/C6 and C3/C5 move downfield and the averaged peaks for the exocyclic carbons C1' and C2'/C3' move upfield at higher temperature as the population of the less stable isomer **2B** is increasing. The averaged chemical shift for the methylene carbon C4, however, moves in the opposite direction to the shifts of the other ring carbons. We attribute this to an additional fast equilibrium of two conformational isomers (hyperconjomers) of the 1-(2-propyl)cyclohexyl cation **2B1** and **2B2** similar as reported for the 1-methylcyclohexyl cation.^{17,18}

Quantum chemical calculations at the DFT, MP2 and MP4 levels of theory reveal that, indeed, three different minimum structures of similar energy (Scheme 3, Table 2) are involved in



equilibrium **2A** $\rightleftharpoons$ **2B** (ΔE ca. ± 1 kcal mol⁻¹, at the MP4/6-31G(d)//MP2/6-31G(d) level).

Two different conformations for the 1-(2-propyl)cyclohexyl cation, **2B1** and **2B2**, and the 2-cyclohexyl-2-propyl cation **2A** are presented in Fig. 3. The μ -hydrido-bridged structure **2TS** was found as a stationary point and characterized as a transition state structure by frequency calculations (NImag = 1) at the B3LYP/6-31G(d) and MP2/6-31G(d) levels of theory (Fig. 3).

The 2-cyclohexyl-2-propyl cation **2A** shows characteristic structural parameters which indicate a CC-hyperconjugative interaction of the formally empty $2p\pi$ -orbital of the C^+ carbon C1' with one of the β -C–C bonds of the cyclohexyl ring. The interacting cyclohexyl C1–C6 bond is perpendicular (89° , at the MP2 level) to the C1–C2'–C3' plane centered on the sp^2 -hybridized C^+ carbon C1', thus the C1–C6 bond is eclipsed to the formal vacant $2p\pi$ orbital at C1'. The bond angle C6–C1–C1' (94.9°) is reduced, the bond angle C5–C1–C1' (121.1°) is increased compared to the C–C–C-bond angles (*ca.* 110°) of the cyclohexyl ring. The C1–C6 bond is elongated (1.626 Å) and the C1'–C1 bond is shortened (1.420 Å) as compared to the other cyclohexyl ring bonds (1.51–1.53 Å).

Both conformational isomers of the 1-(2-propyl)cyclohexyl cation, structures **2B1** and **2B2**, have a chair conformation of the cyclohexyl ring. In **2B1** the axis of the formally empty $2p\pi$ -orbital of the C^+ carbon is in an equatorial-like position, close to parallel to the C2–C3–C4–C6 plane, while in the conformer **2B2**, the empty $2p\pi$ -orbital occupies the axial-like position. The structures **2B1** and **2B2** resemble the two conformations calculated for the 1-methylcyclohexyl cation.¹⁸ In both structures **2B1** and **2B2** one methyl group (C2'H₃) is oriented suitably for CC-hyperconjugation of the vacant C^+ - $2p\pi$ -orbital and the β -C1'–CH₃ bond. The C1'–C2' bond to this methyl group is longer than the bond to the other methyl group (C3'H₃) ($d(C1'–C2') = 1.575$ Å *vs.* $d(C1'–C3') = 1.525$ Å in **2B1**, and $d(C1'–C2') = 1.585$ Å *vs.* $d(C1'–C3') = 1.525$ Å in **2B2**, at the MP2/6–31G(d) level). In the conformational isomer **2B1**, the alignment of the C^+ - $2p\pi$ -orbital and the ring C2–C3 and C5–C6 carbon bonds allows additional stabilization of the positive charge by β -CC hyperconjugation. As a consequence the C2–C3 and C5–C6 bonds are elongated (1.587 Å) and the C1–C2 and C1–C6 bonds are shortened (1.450 Å) compared to the C3–C4 and C4–C5 bonds (1.52 Å). In the isomer **2B2** the two axial β -hydrogens at the methylene carbons C2 and C6 are aligned with the vacant C^+ - $2p\pi$ -orbital, thus involved in β -CH hyperconjugative stabilization of the positive charge. The C2–H and C6–H bonds are slightly longer (1.110 Å, at the MP2/6–31G(d) level) than the average of the other C–H bonds (1.094 Å) of the cyclohexyl ring. Similarly to **1A/1B**, no minimum energy structures with a conformation suitable for β -methine-C–H/ C^+ - 2π -hyperconjugation was found.

The quantum chemical calculations show that the cation structure **2A** and the structures **2B1** and **2B2** (Table 2) are very similar in energy. The small energy differences indicate a very flat energy surface. The order of relative energies calculated for the gas phase structures of **2A** and **2B1/2B2** are not in accord with the experimental results which show that in solution **2A** is preferred in the equilibrium (see Table 2). To estimate whether the effect of solvation^{19,20} could change the relative order of energies for **2A**, **2B1** and **2B2**, self-consistent reaction field (SCRF) calculations with a static isodensity surface polarized continuum model (IPCM)²¹ were performed at B3LYP/6–31G(d) and MP2/6–31G(d) levels. A value of $\epsilon = 30.0$ was used as an input parameter for a solvent relative permittivity. The relative permittivity of the solvent SO₂Cl₂ has been reported ($\epsilon = 9.1$),²² but an excess of SbF₅ would create a more polar medium.¹⁸ The relative order of the energies in solution calculated at the B3LYP/6–31G(d) level of theory is in accord with the experimental findings (Table 1). Structure **2A** in a solution is 1.0 kcal mol^{−1} lower in energy than **2B1** and 0.9 kcal mol^{−1} lower in energy than **2B2**. Applying the solvent model calculation (SCRF = IPCM, $\epsilon = 30.0$) to the cyclopentyl type carbocations **1A** and **1B**, however, does not change the relative order of their stability. The relative stability of **1B** over **1A** is 1.5 kcal mol^{−1} in the gas phase and 1.6 kcal mol^{−1} in the solution model calculation at the B3LYP/6–31G(d) level of theory (Table 2). Thus, for both carbocation systems **1A/1B** and **2A/2B** solvent model calculations at the B3LYP/6–31G(d) level are in qualitative agreement with the experiment.

The different temperature dependence for the cyclohexyl ring carbon of the equilibrating carbocations **2A/2B**, *i.e.* the shielding of the averaged peak for C4 and the deshielding for the averaged peaks for C3/C5, C2/C6 and C1 can be explained taking into account the relative energies and the calculated and scaled chemical shifts of the structures **2A**, **2B1** and **2B2** involved in the equilibrium process (Tables 2 and 3). Raising the temperature, the average shifts of the ring carbons C3/C5, C2/C6 and C1 increase, as the fraction of the lowest energy structure **2A** decreases. The calculated shift for carbon C4 in the more stable **2B** conformer, **2B1** ($\delta_{\text{scaled}} = 30.8$ ppm), is only slightly downfield from the C4 shift in **2A** ($\delta_{\text{scaled}} = 27.5$ ppm), while the shift for carbon C4 in **2B2** ($\delta_{\text{scaled}} = 26.7$ ppm) is slightly upfield from the signal for C4 in **2A**. When the temperature is raised, the fraction of component **2B2** increases more than the fraction of its more stable conformational isomer **2B1** and the net effect is a slight shielding effect for the averaged peak of carbon C4. At higher temperatures the population of **2B1** (C3/C5, $\delta_{\text{scaled}} = 60.7$ ppm) increases relative to **2A** (C3/C5, $\delta_{\text{scaled}} = 29.7$ ppm) causing deshielding of the peak for C3/C5, while an increase of the population of **2B2** (C3/C5, $\delta_{\text{scaled}} = 27.6$ ppm) has the opposite effect, and causes an upfield shift of the C3/C5 peak. Even though the resulting averaged peak for the C3/C5 carbons is shifted downfield at higher temperatures, as expected due to increased charge at these positions, the change in chemical shift is very small ($\Delta\delta = 0.26$ ppm for $\Delta T = 40$ K). This is due to the large chemical shift difference between the C3/C5 signals of **2A** and **2B1**, ($\Delta\delta = 31$ ppm), and the small chemical shift difference for the C3/C5 signals in **2A** and **2B2** ($\Delta\delta = 2.1$ ppm).

The three carbocation structures **2A**, **2B1** and **2B2** involved in the equilibrium (Scheme 3) have the relative population: [**2A**] : [**2B1**] : [**2B2**] = 1 : K_1 : K_2 . The averaged chemical shift of any carbon atom is given by eqn. (4),

$$\delta = \frac{\delta(2A) + K_1\delta(2B1) + K_2\delta(2B2)}{1 + K_1 + K_2} \quad (4)$$

where $\delta(2A)$, $\delta(2B1)$, and $\delta(2B2)$ represent the chemical shift of the carbon in the structures **2A**, **2B1** and **2B2**, respectively. Attempts to fit the corresponding equilibrium constants K_1 and K_2 (Scheme 3) led to considerable error, mainly due to the fact that both the relative energies ($\Delta E = 0.4$ kcal mol^{−1} at the B3LYP/6–31G(d) level) and the chemical shifts (Table 3) are very similar in both conformers **2B1** and **2B2**. The equilibrium constants given in Table 1 represent the sum of the two equilibrium constants $K_{\text{obs}} = K_1 + K_2$ and were calculated using eqn. (3) with the approximation that $K_1 = K_2$. If K_1 is equal to K_2 , the equilibrium constant for both processes **2A** $\rightleftharpoons$ **2B1** and **2A** $\rightleftharpoons$ **2B2** is $K_{\text{obs}}/2$. The corresponding thermodynamic parameters obtained from a van't Hoff plot ($1/T$ plotted *vs.* $\ln(K_{\text{obs}}/2)$) are: $\Delta H = 2.5 \pm 0.2$ kcal mol^{−1}, and $\Delta S = -7.0 \pm 0.8$ cal mol^{−1} K^{−1} indicating that the structure **2A** is energetically favoured over **2B1** and **2B2** by about 2 kcal mol^{−1} which is in fair agreement with the calculated values (Table 2).

The ¹³C NMR spectra of cation **2** show line broadening at lower temperatures. The limits of the rate constants for the equilibration process **2A** $\rightleftharpoons$ (**2B1** + **2B2**)/2 were estimated to be similar to those of cation **1A/1B**, $4.0 \times 10^4 < k_{\text{obs}} < 1.8 \times 10^6$ (−122 °C). The rate constant k_{obs} refers to the sum of the overall rate constants for both processes **2A** $\rightarrow$ **2B1** (k_1) and **2A** $\rightarrow$ **2B2**, (k_2). The single rate constants k_1 and k_2 are approximately $k_{\text{obs}}/2$. This limits the free energy of activation ($\Delta G^\ddagger$) at −122 °C for both processes to the range $4.3 < \Delta G^\ddagger < 5.5$ kcal mol^{−1}. This is somewhat lower than the calculated energy difference between **2A** and **2TS** (7.8 and 7.6 kcal mol^{−1} at the B3LYP/6–31G(d) and MP4/6–31G(d)/MP2/6–31G(d) levels of theory, respectively).

Equilibrium isotope effects in the 1-(2-propyl)cyclopentyl-2-cyclopentyl-2-propyl cation and the 1-(2-propyl)cyclohexyl-2-cyclohexyl-2-propyl cation

The measurement of equilibrium isotope effects (EIE) on NMR spectra — the isotopic perturbation method — is a well established tool for the investigation of rapidly rearranging molecules which show averaged NMR signals.⁷ The deuterium EIE (K_H/K_D) is mainly vibrational in origin, caused by zero-point energy differences between exchanging sites. The heavier isotope deuterium prefers the site with the stiffer bond, where stretching and bending force constants are larger.²³

Therefore in deuterated carbocations where H and D are exchanged between a β -site loose-bonding situation and a stiffer bonding at a more remote site, the equilibrium is shifted towards the isotopomer with deuterium at the remote site.

In the methyl hexadeuterated carbocation **1A-d₆**/**1B-d₆** and **2A-d₆**/**2B-d₆** the equilibria will be shifted towards the cation isomers **1B-d₆** and **2B-d₆**, respectively. If the β -methylene positions, *i.e.* the ring-CH₂ groups are deuterated, as in **1-d₄** and **2-d₄**, the corresponding equilibria will be shifted towards the 2-propyl cation isomers **1A** and **2A**, respectively.

In the ¹³C spectrum of a mixture of **1A**/**1B** and **1A-d₄**/**1B-d₄** (Fig. 4a) the splitting between the C1 and C1' peaks are smaller for **1A-d₄**/**1B-d₄** ($\Delta_D = 202.8$ ppm at -130 °C) than for **1A**/**1B** ($\Delta_H = 271.9$ ppm at -130 °C), while in the ¹³C NMR spectrum of a mixture of **1A**/**1B** and **1A-d₆**/**1B-d₆** (Fig. 4b) the splitting

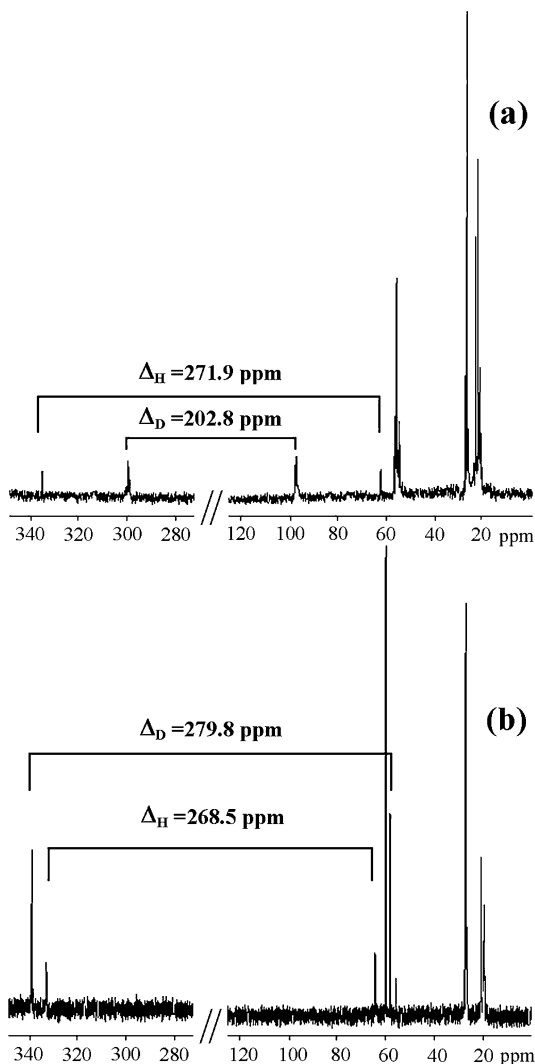


Fig. 4 ¹³C NMR spectra (100 MHz) in SO₂ClF–SO₂F₂ solution: (a) the mixture of **1A**/**1B** and tetraduterated isotopomers **1A-d₄**/**1B-d₄** at -130 °C; (b) the mixture of **1A**/**1B** and hexadeuterated isotopomers **1A-d₆**/**1B-d₆** at -127 °C.

between the C1 and C1' peaks are larger ($\Delta_D = 279.8$ ppm at -127 °C) for **1A-d₆**/**1B-d₆** than for **1A**/**1B** ($\Delta_H = 268.5$ ppm at -127 °C).

These results are consistent only if C1 corresponds to the signal at lower field and C1' to the signal at higher field, *i.e.* the positive charge in the major isomer is located on C1. The EIE in the C2',C3'-hexadeuterated cation **1A-d₆**/**1B-d₆** shifts the non-degenerate equilibrium between the isomeric structures **A** and **B** further towards the cyclopentyl cation structure **B**. This results in a further downfield shift of C1 and corresponding upfield shift of C1', thus a larger splitting, as compared to the non-deuterated cation **1A**/**1B**, is observed for the C1–C1' peak pair. The equilibrium isotope effect for the equilibrium **1B** $\rightleftharpoons$ **1A** is $K_{eq}(H)/K_{eq}(D) = 0.478$ at -112 °C, *i.e.* the average isotope effect per deuterium is 0.88 ($K_H/K_D = 1.14$ per deuterium). The EIE in the C2,C3-tetraduterated cation **1A-d₄**/**1B-d₄** favours the 2-propyl cation isomer **1A-d₄**. The concentration of isomer **1B-d₄** is decreased as compared to non-deuterated equilibrium **1A**/**1B**. This results in an upfield shift of the C1 peak and a downfield shift of the C1' peak, and in the ¹³C spectrum a reduced splitting for the C1 and C1' peak pair is observed for **1A-d₄**/**1B-d₄** compared to the splitting for **1A**/**1B**. The $K_{eq}(H)/K_{eq}(D) = 4.54$ at -116 °C (average EIE per deuterium $K_H/K_D = 1.45$). The larger average EIE per deuterium in the methylene deuterated cation **1A-d₄**/**1B-d₄** as compared to the methyl deuterated cation **1A-d₆**/**1B-d₆** is in accord with earlier conclusions comparing kinetic isotope effects in solvolysis reactions of β -methylene and β -methyl deuterated compounds.²⁴

The changes in the ¹³C NMR spectra of the deuterated isomers **2A-d₄**/**2B-d₄** and **2A-d₆**/**2B-d₆** are opposite compared to those of cations **1A-d₄**/**1B-d₄** and **1A-d₆**/**1B-d₆**. For the C2,C6 methylene deuterated carbocation **2A-d₄**/**2B-d₄** (Fig. 5a) the shift difference between C1 and C1' is larger ($\Delta_D = 246.4$ ppm at -122 °C) than for **2A**/**2B** ($\Delta_H = 234.5$ ppm at -122 °C).

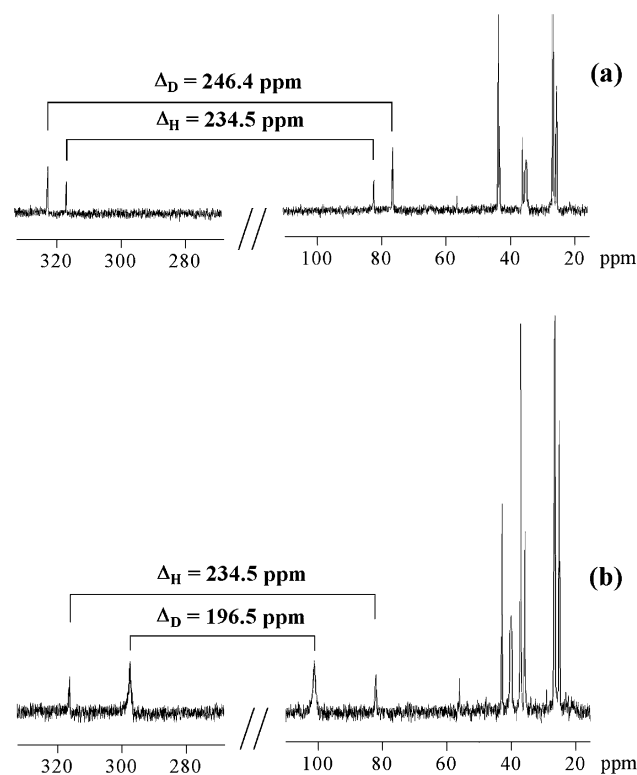
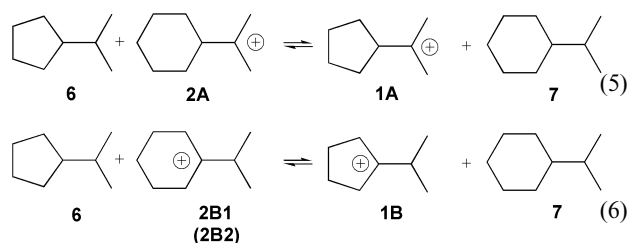


Fig. 5 ¹³C NMR (100 MHz) spectra in SO₂ClF–SO₂F₂ solution: (a) the mixture of **2A**/**2B** and tetraduterated isotopomers **2A-d₄**/**2B-d₄** at -122 °C; (b) the mixture of **2A**/**2B** and hexadeuterated isotopomer **2A-d₆**/**2B-d₆** at -122 °C.

For cation **2A-d₆**/**2B-d₆** the chemical shift difference (Fig. 5b) between C1 and C1' is smaller ($\Delta_D = 196.5$ ppm at -122 °C) than for the non-deuterated cation **2A/2B** ($\Delta_H = 234.5$ ppm at -122 °C). $K_{eq}(D) < K_{eq}(H)$, the isotope effect favours structures of type **B**, thus the concentration of isomer **A** is lower compared to the non-deuterated cation. This shows that the peak at higher field arises from carbon C1', while the peak at lower field is due to carbon C1. The EIE favours structures **B**, therefore C1' is shifted upfield and C1 is shifted downfield in **2A-d₆**/**2B-d₆** and the splitting of the peak pair C1/C1' is reduced as compared to the nondeuterated cation **2A/2B**. The EIE for the **2A** $\rightleftharpoons$ (**2B1/2B2**) equilibria both favour the 2-propyl cation isomer **2A**. The EIE for the **2A-d₆**/**2B-d₆** is $K_{eq}(H)/K_{eq}(D) = 0.478$ at -113 °C; the average isotope effect per deuterium of $K_H/K_D = 1.13$, while for the tetradeuterated carbocation **2A-d₄**/**2B-d₄**, the EIE is $K_{eq}(H)/K_{eq}(D) = 9.2$ at -99 °C, which corresponds to $K_H/K_D = 1.74$ per deuterium.

Isodesmic reactions

The relative stabilities of the carbocation structures **1A/1B** and **2A/2B1/2B2** were evaluated from the isodesmic eqns. (5) and (6).



The relative energies calculated for the 2-cyclohexyl and 2-cyclopentyl substituted 2-propyl cations **2A** and **1A** and the corresponding hydrocarbons isopropylcyclopentane **6** and isopropylcyclohexane **7** (eqn. 5) indicate only a slight preference of the 2-cyclopentyl-2-propyl cation **1A** (ΔE_1 ($\epsilon = 1.0$) = -0.46 kcal mol $^{-1}$, ΔE_1 ($\epsilon = 30.0$) = -1.56 kcal mol $^{-1}$) compared to the cyclohexyl substituted cation **2A**. Thus the relative stability of the 2-(2-cycloalkyl)propyl cations **1A** and **2A** is not very sensitive to the ring size of the cycloalkyl substituent. A larger effect of the ring size, however, is found for 1-(2-propyl)-cycloalkyl cations **1B** and **2B1**, **2B2** (eqn. 6). The 2-propyl substituted cyclopentyl cation **1B** is thermodynamically more stable than the 2-propyl substituted cyclohexyl-cation structures **2B1** and **2B2** (ΔE_2 ($\epsilon = 1.0$) = -1.59 kcal mol $^{-1}$, ΔE_2 ($\epsilon = 30.0$) = -4.16 kcal mol $^{-1}$ and ΔE_2 ($\epsilon = 1.0$) = -1.76 kcal mol $^{-1}$, ΔE_2 ($\epsilon = 30.0$) = -4.06 kcal mol $^{-1}$, respectively).

Conclusions

The nondegenerate 1,2-hydride shifts in (2-propyl)cyclopentyl and (2-propyl)cyclohexyl carbocations **1A/1B** and **2A/2B** were measured by dynamic ^{13}C NMR spectroscopy in superacid solution. The direction of these equilibria depends on the size of the ring. The cyclopentyl cation **1B** is more stable than the 2-propyl cation isomer **1A**, but the 2-propyl cation **2A** is more stable than the two hyperconjugative isomers (hyperconjomers) **2B1** and **2B2** of the cyclohexyl cation **2B**. The results were confirmed by NMR spectroscopic measurements of deuterium equilibrium isotope effects (EIE) and rationalized by quantum chemical calculations of relative energies, geometries and NMR chemical shifts.

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References

- G. A. Olah and D. J. Donovan, *J. Am. Chem. Soc.*, 1977, **99**, 5026.
- D. A. Forsyth and Y. Pan, *Tetrahedron Lett.*, 1985, **26**, 4997.
- N. E. Okazawa and T. S. Sorensen, *Can. J. Chem.*, 1982, **60**, 2180.
- M. Saunders, *et al.* unpublished results, Yale University; J. R. Lloyd, PhD Thesis, Yale University, 1978.
- Only the labeled structures of the more stable structures in the equilibrium are presented.
- M. Saunders and P. Vogel, *J. Am. Chem. Soc.*, 1971, **93**, 2559; D. Lenoir and H.-U. Siehl, *Houben-Weyl: Methoden der Organischen Chemie*, vol. E19C, Georg Thieme, Stuttgart, 1990; p. 26.
- M. Saunders and H. A. Jiménez-Vazquez, *Chem. Rev.*, 1991, **91**, 375; H.-U. Siehl, *Adv. Phys. Org. Chem.*, 1987, **23**, 63.
- Δ represents the chemical shift difference of the interchanging sites in the slow exchange limiting structures, and δ represents the experimental splitting of the interchanging sites.
- V. Vrček, O. Kronja and H.-U. Siehl, *J. Chem. Soc., Perkin Trans. 2*, 1999, 1317.
- Gaussian 98, Revision A.9, M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, V. G. Zakrzewski, J. A. Montgomery, Jr., R. E. Stratmann, J. C. Burant, S. Dapprich, J. M. Millam, A. D. Daniels, K. N. Kudin, M. C. Strain, O. Farkas, J. Tomasi, V. Barone, M. Cossi, R. Cammi, B. Mennucci, C. Pomelli, C. Adamo, S. Clifford, J. Ochterski, G. A. Petersson, P. Y. Ayala, Q. Cui, K. Morokuma, D. K. Malick, A. D. Rabuck, K. Raghavachari, J. B. Foresman, J. Cioslowski, J. V. Ortiz, A. G. Baboul, B. B. Stefanov, G. Liu, A. Liashenko, P. Piskorz, I. Komaromi, R. Gomperts, R. L. Martin, D. J. Fox, T. Keith, M. A. Al-Laham, C. Y. Peng, A. Nanayakkara, M. Challacombe, P. M. W. Gill, B. Johnson, W. Chen, M. W. Wong, J. L. Andres, C. Gonzalez, M. Head-Gordon, E. S. Replogle and J. A. Pople, Gaussian, Inc., Pittsburgh PA, 1998.
- V. Vrček, H.-U. Siehl and O. Kronja, *J. Phys. Org. Chem.*, 2000, **13**, 616.
- J. H. Botkin, D. A. Forsyth and D. J. Sardella, *J. Am. Chem. Soc.*, 1986, **108**, 2797.
- P. v. R. Schleyer, C. Maerker, P. Buzek and S. Sieber, in *Stable Carbocation Chemistry*, Eds. G. K. S. Prakash and P. v. R. Schleyer, John Wiley & Sons, New York, 1997, ch. 2.
- J. Sandström, *Dynamic NMR spectroscopy*, Academic Press, London, 1982, p. 77.
- M. Saunders and M. R. Kates, *J. Am. Chem. Soc.*, 1978, **100**, 7082.
- The lower limit: $k > \pi/[v_{1A}(C1) - v_{1B}(C1)]$, where $v_{1A}(C1)$ and $v_{1B}(C1)$ represent the calculated absorption frequencies of C1 in **1A** and **1B** respectively; The upper limit: $k < \pi[v_{1A}(C2') - v_{1B}(C2')]^2/2\omega$ where $v_{1A}(C2')$ and $v_{1B}(C2')$ represent the frequencies of the C2', C3'-methyl carbons sites in **1A** and **1B** and ω is the line broadening (7 Hz at -127 °C) observed for the methyl carbons C2', C3'.
- R. P. Kirchen, K. Ranganayakulu and T. S. Sorensen, *J. Am. Chem. Soc.*, 1987, **109**, 7811.
- A. Rauk, T. S. Sorensen and P. v. R. Schleyer, *J. Chem. Soc., Perkin Trans. 2*, 2001, 869; A. Rauk, T. S. Sorensen, C. Maerker, J. W. de M. Carnerio, S. Sieber and P. v. R. Schleyer, *J. Am. Chem. Soc.*, 1996, **118**, 3761.
- W. L. Jorgensen, *J. Am. Chem. Soc.*, 1977, **99**, 280; C. Jenson and W. L. Jorgensen, *J. Am. Chem. Soc.*, 1997, **119**, 10846.
- P. v. R. Schleyer, J. W. de M. Carnerio, W. Koch and D. A. Forsyth, *J. Am. Chem. Soc.*, 1991, **113**, 3990; P. R. Schreiner, D. L. Severance, W. L. Jorgensen, P. v. R. Schleyer and H. F. Schaefer III, *J. Am. Chem. Soc.*, 1995, **117**, 2663.
- J. B. Foresman, T. A. Keith, K. B. Wiberg, J. Snoonian and M. J. Frisch, *J. Phys. Chem.*, 1996, **100**, 16098.
- Handbook of Chemistry and Physics*, 80th edn., Ed. D. R. Lide, CRC Press, New York, 1999–2000.
- For general treatment of isotope effects see: J. Biegeleisen and M. Wolfsberg, *Adv. Chem. Phys.*, 1958, **1**, 15; L. Melander, *Isotope Effects on Reaction Rates*, Ronald Press, New York, 1960; *Isotope Effects in Chemical Reactions*, Eds. C. I. Collins and N. S. Bowman, ACS Monograph 167, Van Norstrand Reinhold, New York, 1970; M. Wolfsberg, *Acc. Chem. Res.*, 1972, **7**, 225; L. Melander and W. H. Saunders, Jr., *Reaction Rates of Isotopic Molecules*, Wiley, New York, 1980.
- V. J. Shiner, Jr., *J. Am. Chem. Soc.*, 1961, **83**, 240; V. J. Shiner, Jr. and J. S. Humphrey, Jr., *J. Am. Chem. Soc.*, 1963, **85**, 2146.