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Measurement of the cross section for electroweak production of $Z\gamma$ in association with two jets and constraints on anomalous quartic gauge couplings in proton-proton collisions at $\sqrt{s} = 8\,\text{TeV}$

The CMS Collaboration[*]

Abstract

A measurement is presented of the cross section for the electroweak production of a Z boson and a photon in association with two jets in proton-proton collisions at $\sqrt{s} = 8\,\text{TeV}$. The Z bosons are identified through their decays to electron or muon pairs. The measurement is based on data collected with the CMS detector corresponding to an integrated luminosity of $19.7\,\text{fb}^{-1}$. The electroweak contribution has a significance of 3.0 standard deviations, and the measured fiducial cross section is $1.86^{+0.90}_{-0.75}\,(\text{stat})^{+0.34}_{-0.26}\,(\text{syst}) \pm 0.05\,(\text{lumi})\,\text{fb}$, while the summed electroweak and quantum chromodynamic total cross section in the same region is observed to be $5.94^{+1.53}_{-1.35}\,(\text{stat})^{+0.43}_{-0.37}\,(\text{syst}) \pm 0.13\,(\text{lumi})\,\text{fb}$. Both measurements are consistent with the leading-order standard model predictions. Limits on anomalous quartic gauge couplings are set based on the $Z\gamma$ mass distribution.

*See Appendix A for the list of collaboration members

1 Introduction

With the discovery of the Higgs boson at the CERN LHC [1, 2], the standard model (SM) became a great success. The high energy and luminosity of the LHC provides the opportunity to observe many processes that are predicted by the SM, including electroweak production of multiple gauge bosons (WVγ [3], V$\gamma\gamma$ [4–6]), vector boson scattering (VBS) (same charge W$^\pm$W$^\pm$ scattering [7–9], $\gamma\gamma \rightarrow$ W$^+$W$^-$ [10], EW Wγjj [11], W$^\pm$Z [12]), and vector boson fusion (VBF) (EW W(Z)jj [13–16]). Same charge W$^\pm$W$^\pm$ scattering has been observed by ATLAS, and the exclusive $\gamma\gamma \rightarrow$ W$^+$W$^-$ process by CMS, both with significances larger than 3 standard deviations. The triboson final state Z$\gamma\gamma$ has been observed by ATLAS and CMS with a significance larger than 5 standard deviations. The EW production of a Z boson (decaying into two oppositely-charged leptons), a photon, and two jets (henceforth denoted Zγjj) has never been studied before, and is the subject of this paper. While the cross section for quantum chromodynamic (QCD) induced Zγjj production is orders of magnitude larger than the one for EW production, the latter can be used to perform important tests of the SM, and to search for contributions from physics beyond the SM that could manifest themselves as anomalous trilinear or quartic gauge boson couplings (aTGC or aQGC [3–7, 9, 10, 12**?**]).

This letter presents a measurement of the associated EW production of Zγjj, using the 8 TeV proton-proton collision data recorded by the CMS detector. The major processes contributing to EW Zγjj production are represented by the Feynman diagrams in Fig. 1. They are (a) bremsstrahlung, (b) multiperipheral (or non-resonant) production, (c,d) VBF with either two trilinear gauge boson couplings (TGC), or (e) VBS with quartic gauge boson couplings (QGC). The VBS processes are particularly interesting because they involve QGCs (e.g. WWZγ). It is not possible, however, to isolate the QGC processes from the other contributions, such as the double TGC processes that are topologically similar. The interference of the VBS diagrams ensures unitarity of the VBS cross section in the SM at high energy. We present measurements of the combined cross sections for all EW processes that result in the Zγjj final state. The main background source is Zγjj production where the associated jets are produced through QCD-induced processes (such as the Feynman diagram given in Fig. 1(f)). Other backgrounds include jets or leptons misidentified as photons, diboson processes in which a W or Z boson decays into two jets and the photon originates from initial or final-state radiation, and contributions from top quark pairs and single top quark production.

2 The CMS detector

The central feature of the CMS apparatus is a superconducting solenoid of 6 m internal diameter, providing a magnetic field of 3.8 T. Within the solenoid volume are a silicon pixel and strip tracker, a lead tungstate crystal electromagnetic calorimeter (ECAL), and a brass and scintillator hadron calorimeter (HCAL), each composed of a barrel and two endcap sections. Forward calorimeters extend the pseudorapidity, η, coverage provided by the barrel and endcap detectors. Muons are measured in gas-ionization detectors embedded in the steel flux-return yoke outside the solenoid.

The particle-flow (PF) event algorithm [17, 18] reconstructs and identifies each individual particle with an optimized combination of information from the various elements of the CMS detector. The energy of photons is directly obtained from the ECAL measurement, corrected for zero-suppression effects. The energy of electrons is determined from a combination of the electron momentum at the primary interaction vertex as determined by the tracker, the energy of the corresponding ECAL cluster, and the energy sum of all bremsstrahlung photons spa-
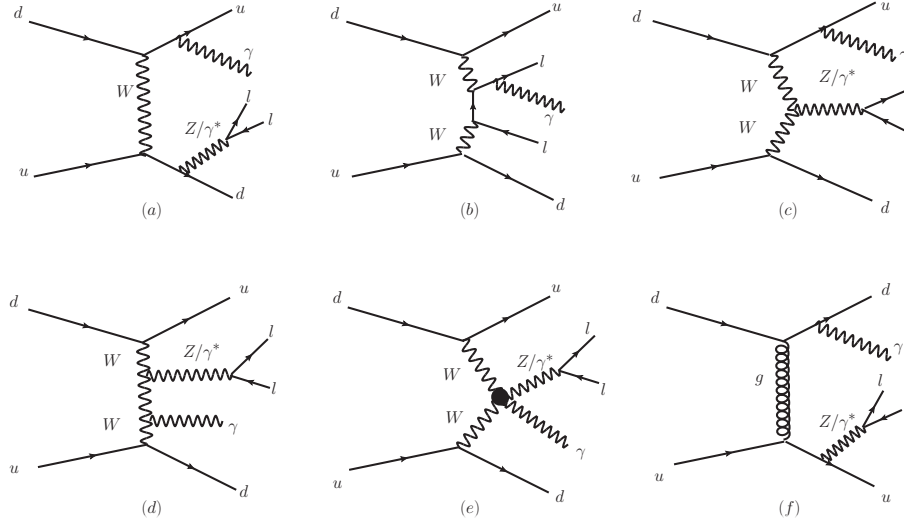
Figure 1: Representative diagrams for EW Zγjj production at the LHC: (a) bremsstrahlung, (b) multiperipheral, (c,d) VBF with TGC, (e) VBS including QGC, and (f) Example diagram for the QCD Zγjj production.

tially compatible with originating from the electron track. The energy of muons is obtained from the curvature of the corresponding track. The energy of charged hadrons is determined from a combination of their momentum measured in the tracker and the matching ECAL and HCAL energy deposits, corrected for zero-suppression effects and for the response function of the calorimeters to hadronic showers. Finally, the energy of neutral hadrons is obtained from the corresponding corrected ECAL and HCAL energy.

In the barrel section of the ECAL, an energy resolution of about 1% is achieved for unconverted or late-converting photons in the tens of GeV energy range. The resolution for other photons in the barrel section is about 1.3% up to $|\eta| = 1$, rising to about 2.5% at $|\eta| = 1.4$. In the endcaps, the resolution for unconverted or late-converting photons is about 2.5%, and the resolution for the remaining photons in the endcap is between 3% and 4% [19]. When combining information from the entire detector, the jet energy resolution is typically 15% at 10 GeV, 8% at 100 GeV, and 4% at 1 TeV.

Muons are measured in the range of $|\eta| < 2.4$, with detection planes utilizing three technologies: drift tubes, cathode strip chambers, and resistive-plate chambers. Matching muons to tracks measured in the silicon tracker results in a p_T resolution for muons with $20 < p_T < 100$ GeV of 1.3–2.0% in the barrel and better than 6% in the endcaps.

The electron momentum is estimated by combining the energy measurement in the ECAL with the momentum measurement in the tracker. The momentum resolution for electrons with transverse momentum $p_T \approx 45$ GeV from Z $\to$ ee decays ranges from 1.7% for nonshowering electrons in the barrel region to 4.5% for showering electrons in the endcaps. The dielectron mass resolution for Z $\to$ ee decays is 1.9% when both electrons are in the ECAL barrel, and 2.9% when both electrons are in the endcaps.

A more detailed description of the CMS detector, together with a definition of the coordinate system used and the relevant kinematic variables, can be found in Ref. [20].

3 Event reconstruction and selection

Candidate events are selected online with triggers that require two muons or electrons, where the leading and subleading leptons have $p_T > 17$ and $8\,\text{GeV}$ respectively, with $|\eta| < 2.4$ (muons) or $|\eta| < 2.5$ (electrons). The overall trigger efficiency is about 94% and 90% for muons and electrons, respectively, with a small dependence on p_T and η.

Muons are reconstructed with a global fit using both the inner tracking system and the muon spectrometer. An isolation requirement is applied in order to suppress the background from multijet events [21, 22]. Electron candidates are reconstructed by matching energy deposits in the ECAL with reconstructed tracks; they must pass stringent quality criteria and an isolation requirement [23]. Charged leptons must originate from the primary vertex, which is defined as the vertex whose tracks have the highest sum of p_T^2. We require that each event has exactly two oppositely charged muons (electrons) with $p_T > 20\,\text{GeV}$ and $|\eta| < 2.4\,(2.5)$ and that the invariant mass of the dilepton system must satisfy $70 < M_{\ell\ell} < 110\,\text{GeV}$. The selection efficiencies for leptons are measured using the tag-and-probe method [24] and are approximately 96% for the muons [25] and 80% for the electrons [21].

Photon candidates are reconstructed from energy deposits in the ECAL with no associated track. Quality selection criteria [19] are applied to the reconstructed photons to suppress the background from hadrons misidentified as photons. The observables used in the photon selection are: (1) PF-based isolation variables that are corrected for the contribution from additional proton-proton collisions in the same bunch crossing (pileup); (2) a small ratio of hadronic energy in the HCAL to electromagnetic energy in the ECAL matched in (η, ϕ) (where ϕ is azimuthal angle in radians); (3) the transverse width of the electromagnetic shower along the η direction [19]; and (4) an electron track veto. We consider only photons in the ECAL barrel region ($|\eta| < 1.44$) with $p_T > 25\,\text{GeV}$. Events with the photon candidate in one of the endcaps ($|\eta| > 1.57$) are excluded from the selection because their signal purity is lower and systematic uncertainties are large.

Hadronic jets are formed from the particles reconstructed by the PF algorithm, using the FAST-JET software package [26] and the anti-k_T jet clustering algorithm [27] with a distance parameter of 0.5. To reduce the contamination from pileup, charged PF candidates in the tracker acceptance region $|\eta| < 2.4$, are excluded from the jet clustering procedure if associated with pileup vertices. The contribution of neutral particles from pileup events to the jet energy is taken into account by means of a correction based on the projected area of the jet on the front face of the calorimeter. Jet energy corrections are derived from a measurement of the p_T balance in dijet and photon+jet events in data [28]. Further residual corrections as functions of p_T and η are applied to the data to correct for the small differences between data and simulation. Additional quality criteria are applied to the jets in order to remove spurious jet-like features originating from isolated noise patterns in the calorimeters or in the tracker [29]. The two jets with the highest p_T are tagged as the signal jets and are required to have $p_T > 30\,\text{GeV}$ and $|\eta| < 4.7$. Since we are primarily interested in the VBS topologies, we require that the invariant mass of the two jets, $M_{jj} > 150\,\text{GeV}$.

Table 1 presents a summary of the three different section criteria that are used for (1) the SM EW signal search, (2) the SM fiducial cross section measurement, and (3) the aQGC searches. The criteria isolate events consistent with the VBS topology of two high-energy scattered jets

separated by a large rapidity gap. The cross section measurement adds two variables sensitive to the VBS process: $|y_{Z\gamma} - (y_{j1} + y_{j2})/2|$, which ensures the $Z\gamma$ systems is located between the scattered jets in eta; and $\Delta\phi_{Z\gamma,jj}$, which requires the $Z\gamma$ system transverse momentum is consistent with recoiling against the transverse momentum of the two combined jets. The fiducial cross section criteria constrain the VBS topology with only basic kinematic cuts that define the acceptance of the CMS detector and a simple two dimensional requirement on the rapidity separation and invariant mass of the jets. A tight p_T^γ selection is applied to reach a higher expected significance in a search for a possible aQGC signal in the EW $Z\gamma jj$ process.

Table 1: Summary of the three different event criteria: (1) selection for the EW signal measurement; (2) the cross section measurement; and (3) the selection for the aQGC search. "j1" and "j2" represent the jets that have the largest and second-largest p_T, "$\ell 1$" and "$\ell 2$" denote the lepton and antilepton from the decay of the Z boson, y is the rapidity, $\Delta\phi_{Z\gamma,jj}$ is the absolute difference between $\phi_{Z\gamma}$ and ϕ_{j1j2}, and the angular separation $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$.

Common selection
$p_T^{j1,j2} > 30\,\text{GeV}$, $\|\eta^{j1,j2}\| < 4.7$
$p_T^{\ell 1,\ell 2} > 20\,\text{GeV}$, $\|\eta^{\ell 1,\ell 2}\| < 2.4$
$\|\eta^\gamma\| < 1.4442$
$M_{jj} > 150\,\text{GeV}$
$70 < M_{\ell\ell} < 110\,\text{GeV}$

EW signal measurement	Fiducial cross section	aQGC search
$p_T^\gamma > 25\,\text{GeV}$	$p_T^\gamma > 20\,\text{GeV}$	$p_T^\gamma > 60\,\text{GeV}$
$\|\Delta\eta_{jj}\| > 1.6$	$\|\Delta\eta_{jj}\| > 2.5$	$\|\Delta\eta_{jj}\| > 2.5$
$\Delta R_{j\ell} > 0.3$, $\Delta R_{jj,\gamma j,\gamma \ell} > 0.5$	$\Delta R_{jj,\gamma j,\gamma \ell,j\ell} > 0.4$	$\Delta R_{j\ell} > 0.3$, $\Delta R_{jj,\gamma j,\gamma \ell} > 0.5$
$\|y_{Z\gamma} - (y_{j1} + y_{j2})/2\| < 1.2$	$M_{jj} > 400\,\text{GeV}$	$M_{jj} > 400\,\text{GeV}$
$\Delta\phi_{Z\gamma,jj} > 2.0$ radians		
$M_{jj} > 400\,\text{GeV}$ with two divided regions		
$400 < M_{jj} < 800\,\text{GeV}$ and $M_{jj} > 800\,\text{GeV}$		

4   Data and simulation

We use data collected with the CMS detector, corresponding to an integrated luminosity of $19.7\,\text{fb}^{-1}$, at proton-proton center-of-mass energy of 8 TeV.

The EW signal, $Z\gamma jj$, at leading-order (LO), and the main background, QCD $Z\gamma$ with 0–3 additional jets, for which the next-to-leading-order (NLO) QCD prediction has been taken from Ref. [30], matched with parton shower based on the so-called "MLM prescription" [31, 32], are simulated using MADGRAPH v5.1.3.30 [33] interfaced with PYTHIA v6.424 [34] for hadronization and showering, using a CTEQ6L1 parton distribution function (PDF) set [35]. The second significant background contribution comes from processes where a jet is misidentified as a photon (fake photon), and this contribution is estimated from data. Other background contributions come from diboson processes (WW/WZ/ZZ) simulated by PYTHIA, single top processes simulated by POWHEG, and $t\bar{t}\gamma$ simulated using MADGRAPH interfaced with PYTHIA. The next-to-leading-order QCD cross sections are used to normalize these simulated samples, except for $t\bar{t}\gamma$ where an LO prediction is taken.

All the simulated events are processed through a GEANT4 [36] simulation of the CMS detector. The tag-and-probe technique is used to correct for data-Monte Carlo (MC) differences in the

trigger efficiency, as well as the reconstruction and selection efficiencies. Additional proton-proton interactions are superimposed over the hard scattering interaction with the distribution of primary vertices matching that obtained from the collision data.

5 Background modeling

The dominant source of background to the EW signal is QCD $Z\gamma$ + jets production. The shape of this background is taken from MC simulation and the normalization is evaluated from data in a control region, defined as $150 < M_{jj} < 400\,\text{GeV}$, where the signal contribution is below 1%. The simulated MC events correctly reproduce the yield of these events with a correction factor of 1.00 ± 0.22 for the combined $Z \to \mu^+\mu^-$ and $Z \to e^+e^-$ channels. The value is comparable with the NLO QCD K factor from Ref. [30], which is around 1.1 for $M_{jj} < 400\,\text{GeV}$.

The background from fake photons arises mainly from Z+jets events where one jet satisfies the photon ID criteria. The estimation is based on events similar to the ones selected with the baseline selection described in Table 1, except that the photon must fail the tight photon ID and satisfy a looser ID requirement based on the charged isolation variable. This selection ensures that the photon arises from a jet, but still has kinematic properties similar to a genuine photon satisfying the tight photon ID. We select genuine photons using $\sigma_{\eta\eta}$, a photon identification variable that exploits the small lateral extension of the electromagnetic shower [19, 25]. Based on the difference between the $\sigma_{\eta\eta}$ distributions for fake photons and genuine photons, a fit is made to normalize the number of events with fake photons to the number of events with genuine photons and obtain the probability to have a fake photon. The fake photon probability is calculated based on different p_T^γ regions in a manner similar to that described in Ref. [37].

Other backgrounds, including top quark and diboson production processes are estimated from MC simulations and normalized to the integrated luminosity of the data sample. The contribution from these backgrounds is less than 10% after applying the kinematic selection (Section 3) and is negligible once the final EW and aQGC selection criteria (Sections 7 and 8) are applied.

The M_{jj} distributions for the $Z \to \mu^+\mu^-$ and e^+e^- channels after the selection requirements described in Section 3 are shown in Fig. 2. The observed distributions are compared to the combined prediction of the backgrounds and of the EW $Z\gamma jj$ signal.

6 Systematic uncertainties

The systematic uncertainty in the QCD $Z\gamma$+jets background estimation is 22% for both $Z \to \mu^+\mu^-$ and $Z \to e^+e^-$; it is dominated by the large statistical uncertainty in the control region used for normalization. The shape uncertainties that are related to the extrapolation of the normalization factor to the signal region ($M_{jj} > 400\,\text{GeV}$) are determined by varying the renormalization and factorization scales as well as the MLM matching scale [31, 32] up and down by a factor of two. Finally, we combine both the normalization factor uncertainty and the shape uncertainty to obtain the total uncertainty.

The systematic uncertainty in the background estimation from fake photons arises from the variation in the choice of the charged isolation sideband and the $\sigma_{\eta\eta}$ distribution used for estimating the fake photon probability. The total uncertainties in the fake photon background estimation can be found in Table 2. The theoretical uncertainty in the top quark background is 20% [3].

The systematic uncertainties in the estimation of the trigger efficiency, measured using the
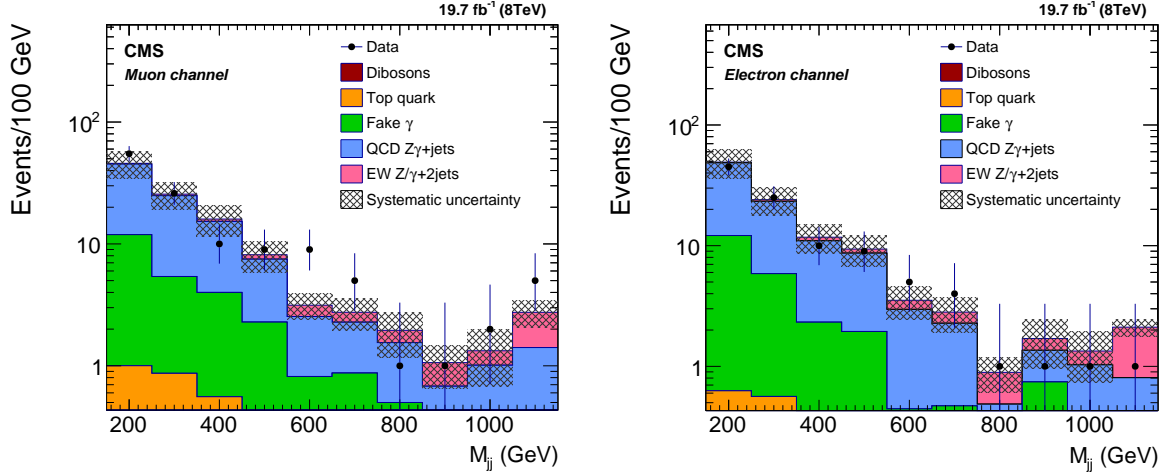
Figure 2: The M_{jj} distributions measured in (left) muon and (right) electron channels. The data (solid symbols with error bars representing the statistical uncertainties) are compared to a data-driven background estimate, combined with MC predictions for the signal contribution. The hashed bands represent the full uncertainty in the predictions, as described in Section 6. The last bin includes overflow events.

tag-and-probe technique, are 1.2% and 1.7% for the $Z \rightarrow \mu^+\mu^-$ and $Z \rightarrow e^+e^-$ channels, respectively. Using similar methods, the systematic uncertainties in the efficiencies for lepton reconstruction and identification in the two channels are 1.9% and 1.0%, respectively. The systematic uncertainty in the jet energy scale and resolution is estimated by varying the jet energy scale and resolution up and down within their p_T- and η-dependent uncertainties [28]. The uncertainty is 14% for $M_{jj} > 400\,\text{GeV}$. Another source of uncertainty is the modelling of the pileup. The inelastic cross section is varied by $\pm 5\%$ in order to evaluate this contribution. The uncertainty in the integrated luminosity is 2.6% [38].

There are also three sources of theoretical uncertainties applied to the signal only. The PDF uncertainty for the signal is estimated with the CT10 [39] PDF set, following the asymmetric Hessian method introduced in Refs. [40, 41]. The scale uncertainty is evaluated by varying the renormalization and factorization scales independently by a factor of two. The magnitude of the interference between QCD and EW $Z\gamma jj$ processes is assigned as systematic uncertainties in the two M_{jj} ranges.

All the systematic uncertainties described are applied to both the signal significance measurement and the aQGC search. They are also propagated to the uncertainty in the measured fiducial cross section, with the exception of the theoretical uncertainty associated with the signal cross section.

All the uncertainties in our analysis are summarized in Table 2.

7   Measurement of the signal significance and fiducial cross section

As shown in Table 1, in addition to the common selection, we apply three further requirements to isolate the EW signal: $|y_{Z\gamma} - (y_{j1} + y_{j2})/2| < 1.2$, $|\Delta\eta_{jj}| > 1.6$, and $\Delta\phi_{Z\gamma,jj} > 2.0$ radians. The selection requirements are chosen by optimizing the expected significance. We apply the CL_s criterion described in Ref. [42, 43] to assess the signal significance, based on the binned

Table 2: Summary of the major uncertainties.

Source	Uncertainty
QCD $Z\gamma$ + jets normalization	22% ($400 < M_{jj} < 800\,\text{GeV}$)
	24% ($M_{jj} > 800\,\text{GeV}$)
Fake photon from jet	15% (20–30 GeV)
(p_T^γ dependent)	22% (30–50 GeV)
	49% (>50 GeV)
Trigger efficiency	1.2% ($Z \to \mu^+\mu^-$), 1.7% ($Z \to e^+e^-$)
Lepton selection efficiency	1.9% ($Z \to \mu^+\mu^-$), 1.0% ($Z \to e^+e^-$)
Jet energy scale and resolution	14% ($M_{jj} > 400\,\text{GeV}$)
$t\bar{t}\gamma$ cross section	20% [3]
Pileup modeling	1.0%
Renormalization/	9.0% ($400 < M_{jj} < 800\,\text{GeV}$), 12% ($M_{jj} > 800\,\text{GeV}$) (SM)
factorization scale (signal)	14% (aQGC)
PDF (signal)	4.2% ($400 < M_{jj} < 800\,\text{GeV}$), 2.4% ($M_{jj} > 800\,\text{GeV}$) (SM)
	4.3% (aQGC)
Interference (signal)	18% ($400 < M_{jj} < 800\,\text{GeV}$), 11% ($M_{jj} > 800\,\text{GeV}$) (SM)
Luminosity	2.6%

M_{jj} distribution, using only the two rightmost bins corresponding to $400 < M_{jj} < 800\,\text{GeV}$ and $M_{jj} > 800\,\text{GeV}$. We consider QCD $Z\gamma jj$ production and events without $Z\gamma$ as background and EW $Z\gamma jj$ production as signal.

Table 3 summarizes the number of events predicted for each process with the number of events observed. For EW $Z\gamma jj$ production, the observations are found to be compatible with expectations in the different channels. By combining both channels, we find evidence for EW $Z\gamma jj$ production with an observed and expected significance of 3.0 and 2.1 standard deviations, respectively. We determine the ratio of the observed signal to that expected from the SM for LO EW $Z\gamma jj$ production as $\hat{\mu} = 1.5^{+0.9}_{-0.6}$ using a binned likelihood fit over the two ranges of the M_{jj} distribution.

Applying the same criteria, we can also measure the significance of the combined EW and QCD $Z\gamma jj$ process. As shown in Table 3, with the two decay channels combined in the search region, of the signal events 7.0 (38.4%) are estimated to come from EW production and the remaining 11.3 from QCD production. As a result, the observed (expected) significance for the combined EW and QCD $Z\gamma jj$ process is 5.7 (5.5) standard deviations.

To determine the cross section for EW $Z\gamma jj$ production we use a fiducial kinematic region based on the acceptance of the CMS detector with a minimal selection on the M_{jj} and $\Delta\eta_{jj}$ variables to select the VBS topology. The fiducial region is defined as described in Table 1. We define the cross section in the fiducial region as $\sigma_f = \sigma_g\,\hat{\mu}\,\alpha_{gf}$ where σ_g is the cross section for generated signal events, $\hat{\mu}$ is the signal strength, and α_{gf} is the acceptance for the generated events in the fiducial region, evaluated through simulation. The fiducial cross section for EW $Z\gamma jj$ production is $1.86^{+0.90}_{-0.75}$ (stat)$^{+0.34}_{-0.26}$ (syst) ± 0.05 (lumi) fb, consistent with the theoretical prediction at LO of 1.27 ± 0.11 (scale) ± 0.05 (PDF) fb calculated using MADGRAPH.

The cross section for all processes that produce the $Z\gamma jj$ final state can be compared to theoretical predictions. The fiducial region studied here lies in a particularly interesting region of phase space because of the substantial contribution to $Z\gamma jj$ from EW production. By restricting the

Table 3: Signal and background yields after the final selection for the SM measurement, for the two bins of $400 < M_{jj} < 800\,\text{GeV}$ (upper) and $M_{jj} > 800\,\text{GeV}$ (lower). Only statistical uncertainties are reported.

$400 < M_{jj} < 800\,\text{GeV}$	muon	electron
Fake photon from jet	3.4 ± 0.8	1.7 ± 0.5
Other background	0.1 ± 0.1	0.1 ± 0.1
QCD $Z\gamma jj$	4.8 ± 0.9	5.0 ± 1.0
EW $Z\gamma jj$	1.7 ± 0.1	1.8 ± 0.1
Total background	8.3 ± 1.2	6.8 ± 1.1
Data	13	8
$M_{jj} > 800\,\text{GeV}$	muon	electron
Fake photon from jet	0.4 ± 0.3	0.1 ± 0.1
Other background	0 ± 0	0 ± 0
QCD $Z\gamma jj$	0.4 ± 0.1	1.1 ± 0.2
EW $Z\gamma jj$	1.8 ± 0.1	1.8 ± 0.1
Total background	0.8 ± 0.3	1.2 ± 0.2
Data	5	2

phase space to the fiducial region for the EW process as defined before, the expected fraction of EW events in the combined sample of EW and QCD signal events is 26%, and the cross section of the combined process is $5.94^{+1.53}_{-1.35}\,\text{(stat)}^{+0.43}_{-0.37}\,\text{(syst)} \pm 0.13\,\text{(lumi)}$ fb, which is consistent with the theoretical prediction at LO calculated using MADGRAPH: $5.05 \pm 1.22\,\text{(scale)} \pm 0.31\,\text{(PDF)}$ fb.

8 Search for anomalous quartic gauge couplings

The effects of any new physics between the TeV and the Planck scale might be significant in the high energy tails of measurements at the LHC and can be parameterized via effective anomalous couplings. With the discovery of the Higgs boson, higher-dimensional operators can be introduced in a linear way [44]:

$$
\begin{aligned}
\mathcal{L}_{\text{aQGC}} =& \frac{f_{M0}}{\Lambda^4}\,\text{Tr}\left[\mathbf{W}_{\mu\nu}\mathbf{W}^{\mu\nu}\right] \times \left[(D_\beta\Phi)^\dagger D^\beta\Phi\right] + \frac{f_{M1}}{\Lambda^4}\,\text{Tr}\left[\mathbf{W}_{\mu\nu}\mathbf{W}^{\nu\beta}\right] \times \left[(D_\beta\Phi)^\dagger D^\mu\Phi\right] \\
&+ \frac{f_{M2}}{\Lambda^4}\left[B_{\mu\nu}B^{\mu\nu}\right] \times \left[(D_\beta\Phi)^\dagger D^\beta\Phi\right] + \frac{f_{M3}}{\Lambda^4}\left[B_{\mu\nu}B^{\nu\beta}\right] \times \left[(D_\beta\Phi)^\dagger D^\mu\Phi\right] \\
&+ \frac{f_{T0}}{\Lambda^4}Tr[\hat{W}_{\mu\nu}\hat{W}^{\mu\nu}] \times Tr[\hat{W}_{\alpha\beta}\hat{W}^{\alpha\beta}] + \frac{f_{T2}}{\Lambda^4}\,Tr[\hat{W}_{\alpha\mu}\hat{W}^{\mu\beta}] \times Tr[\hat{W}_{\beta\nu}\hat{W}^{\nu\alpha}] \\
&+ \frac{f_{T8}}{\Lambda^4}B_{\mu\nu}B^{\mu\nu}B_{\alpha\beta}B^{\alpha\beta} + \frac{f_{T9}}{\Lambda^4}B_{\alpha\mu}B^{\mu\beta}B_{\beta\nu}B^{\nu\alpha},
\end{aligned}
\tag{1}
$$

where $f_{M0,1,2,3}$ and $f_{T0,2,8,9}$ are coefficients of relevant effective operators, and Λ represents the scale of new physics responsible for anomalous couplings. The Lagrangian of the aQGCs is implemented within the MADGRAPH package.

We study the distribution of the mass of the dilepton and photon system, $M_{Z\gamma}$, to search for contributions from aQGCs. The effects of new physics would be seen at higher energy and modify the interference of VBS diagrams. To select the region sensitive to new physics, we require $p_T^\gamma > 60\,\text{GeV}$. The selection for the aQGC analysis is described in Table 1. The $Z\gamma$ mass distribution is shown in Fig. 3, where the last bin includes all events with $M_{Z\gamma} > 420\,\text{GeV}$. Because no significant excess is seen in the $M_{Z\gamma}$ distribution, we use the shape of the $M_{Z\gamma}$ distribution to extract limits on aQGC contributions.
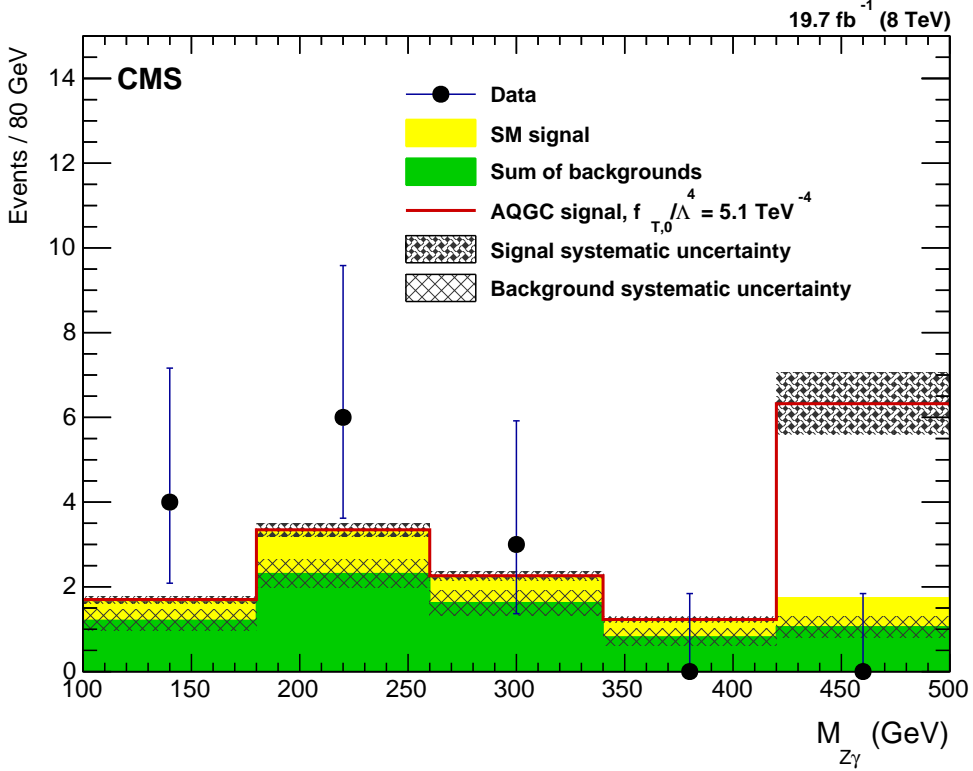
Figure 3: The invariant mass distribution of the $Z\gamma$ system for events that pass the aQGC selection. The highest mass bin includes events with $M_{Z\gamma} > 420\,\text{GeV}$. Error bars represent the statistical uncertainty in the data, while the systematic uncertainties in the aQGC signal and background estimate are shown as hatched bands.

With the parameterization of signals and related systematic uncertainties, for each aQGC parameter, we reweight the SM signal shape to the aQGC shape. The following test statistic is used:

$$t_{\alpha_{\text{test}}} = -2\ln\frac{\mathcal{L}(\alpha_{\text{test}}, \hat{\hat{\boldsymbol{\theta}}})}{\mathcal{L}(\hat{\alpha}, \hat{\boldsymbol{\theta}})}, \tag{2}$$

where the likelihood function ($\mathcal{L}$) is constructed for both lepton channels and combined, using a bin-wise Poisson distribution with profiled nuisance parameters ($\boldsymbol{\theta}$). α_{test} represents the aQGC point being tested. The symbol $\hat{\hat{\boldsymbol{\theta}}}$ represents the values corresponding to the maximum of the likelihood at the point α_{test}, while $\hat{\alpha}$ and $\hat{\boldsymbol{\theta}}$ correspond to the global maximum of the likelihood. This test statistic is assumed to follow a χ^2 distribution [45], from which one can extract limits. Exclusion limits are shown in Table 4. Each coupling parameter is varied over a set of discrete values, keeping the other parameters fixed to zero.

An effective theory is only valid at energies lower than the scale of new physics, and high-dimensional operators with nonzero aQGC values can lead to unitarity violation at sufficiently high energies. For each aQGC listed in Table 4, we checked the stated upper limit against the unitary bound [46] obtained with VBFNLO [47]. In general, we find the limits on all aQGC parameters are set in the unitary unsafe region, except for f_{T9} where the unitarity bound is up to 6 TeV. Form factors can be introduced to unitarize the high energy contribution, however it is difficult to compare results from different experiments and it is not theoretically well motivated. In this study all of the aQGC limits shown are evaluated without a form factor, and can be directly compared to limits set in references [3–7, 9, 10, 12?].

Table 4: Observed and expected shape-based exclusion limits for each aQGC parameter at 95% CL, without a form factor applied.

Observed limits (TeV^{-4})	Expected limits (TeV^{-4})
$-71 < f_{M0}/\Lambda^4 < 75$	$-109 < f_{M0}/\Lambda^4 < 111$
$-190 < f_{M1}/\Lambda^4 < 182$	$-281 < f_{M1}/\Lambda^4 < 280$
$-32 < f_{M2}/\Lambda^4 < 31$	$-47 < f_{M2}/\Lambda^4 < 47$
$-58 < f_{M3}/\Lambda^4 < 59$	$-87 < f_{M3}/\Lambda^4 < 87$
$-3.8 < f_{T0}/\Lambda^4 < 3.4$	$-5.1 < f_{T0}/\Lambda^4 < 5.1$
$-4.4 < f_{T1}/\Lambda^4 < 4.4$	$-6.5 < f_{T1}/\Lambda^4 < 6.5$
$-9.9 < f_{T2}/\Lambda^4 < 9.0$	$-14.0 < f_{T2}/\Lambda^4 < 14.5$
$-1.8 < f_{T8}/\Lambda^4 < 1.8$	$-2.7 < f_{T8}/\Lambda^4 < 2.7$
$-4.0 < f_{T9}/\Lambda^4 < 4.0$	$-6.0 < f_{T9}/\Lambda^4 < 6.0$

9   Conclusions

The measurement of the cross section for the electroweak production of a Z boson and a photon in association with two jets, where the Z boson decays into electron or muon pairs, was presented. The measurement is based on a sample of proton-proton collisions collected with the CMS detector at a center-of-mass energy of 8 TeV, corresponding to an integrated luminosity of 19.7 fb^{-1}. We find evidence for EW Zγjj production with an observed (expected) significance of 3.0 (2.1) standard deviations. The fiducial cross section for EW Zγjj production is measured to be $1.86^{+0.90}_{-0.75}$ (stat)$^{+0.34}_{-0.26}$ (syst) ± 0.05 (lumi) fb, consistent with the theoretical prediction. The fiducial cross section for combined EW and QCD Zγjj production is $5.94^{+1.53}_{-1.35}$ (stat)$^{+0.43}_{-0.37}$ (syst) $\pm$ 0.13 (lumi) fb, which is also consistent with the leading-order theoretical prediction.

In the framework of dimension-eight effective field theory operators, limits on the aQGC parameters $f_{M0,1,2,3}$ and $f_{T0,1,2,8,9}$ are set at 95% confidence level. This is the first constraints on the neutral aQGC parameters f_{T8}.

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